

XI Escola do CBPF
Curso de Pós-Graduação

Notas de Aula

A Física dos Detectores de Partículas

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(web-page: <http://cern.ch/amoraes>)

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Introdução

Programa do Curso:

Aula 1: De Rutherford ao LHC: Desenvolvimento dos detectores ao longo da história da física das partículas elementares. (2^af. 17/07)

Aula 2: Interações das partículas com a matéria. (3^af. 18/07)

Aula 3: Detectando partículas carregadas & neutras. (5^af. 20/07)

Aula 4: Cintiladores: detectando partículas via luminescência. (6^af. 21/07)

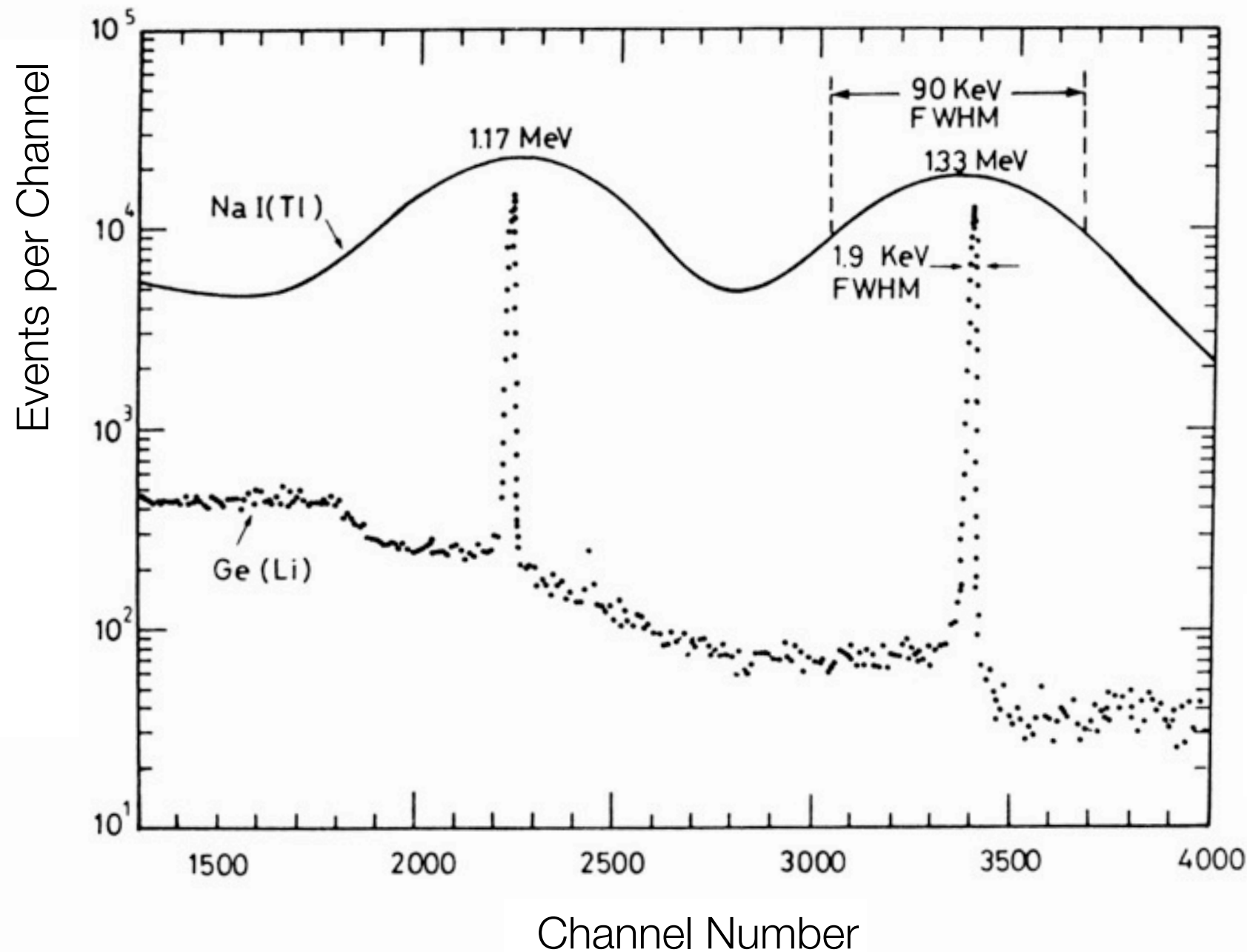
Aula 5: Detectores de semicondutores: medidas de alta precisão. (2^af. 24/07)

Aula 6: Detectores de gás: medindo partículas em grandes volumes. (3^af. 25/07)

Aula 7: Calorímetros: eletromagnéticos & hadrônicos. (5^af. 27/07)

Aula 8: Exemplos de aplicações dos detectores em várias áreas. (6^af. 28/07)

Cintilador vs Semicondutor



Comparison
of two ^{60}Co γ -spectra
one measured with NaI(Tl)
and With Ge(Li) Detector

Aula 6

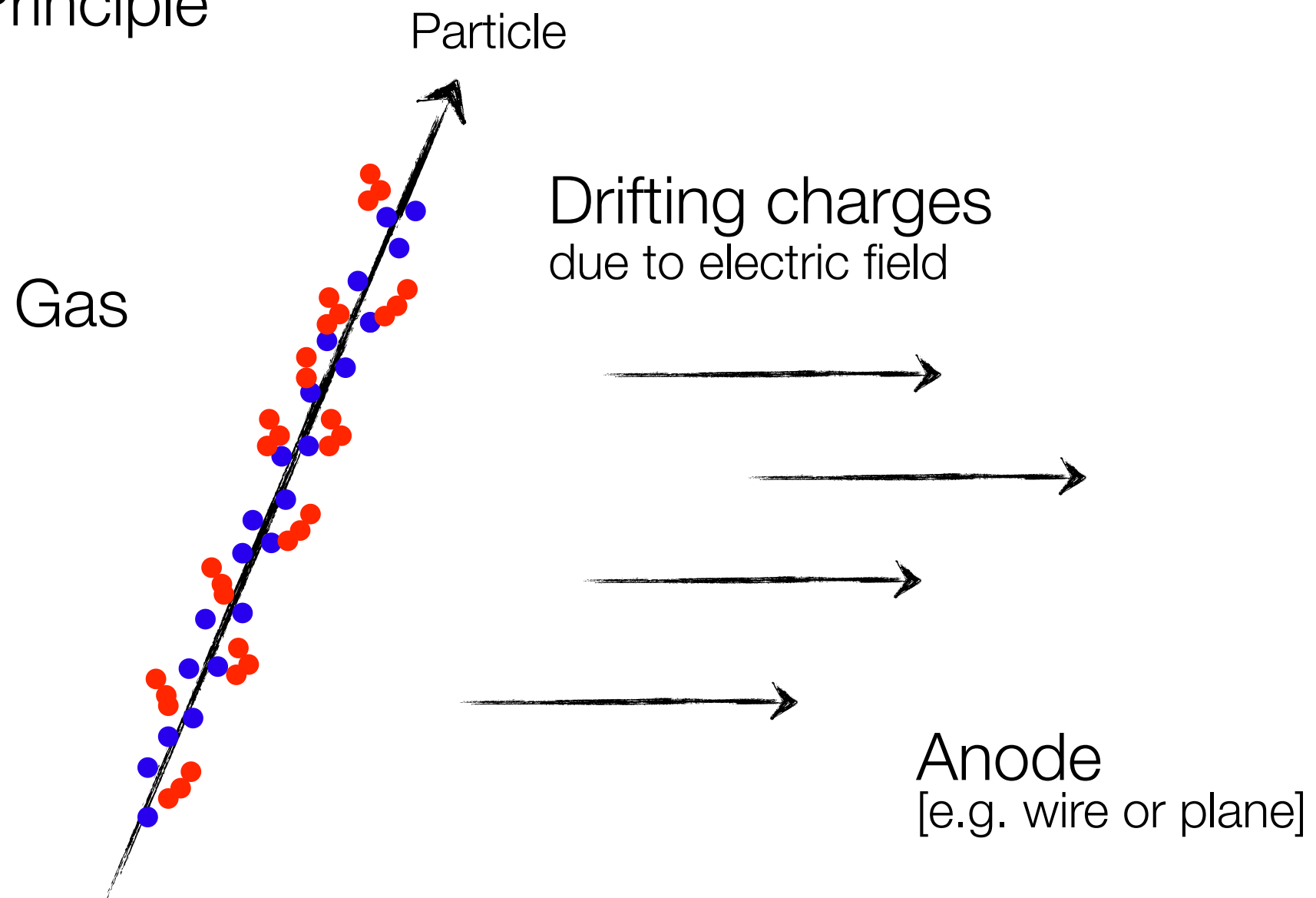
Detectores de gás: medindo partículas em grandes volumes



<http://cern.ch/amoraes>

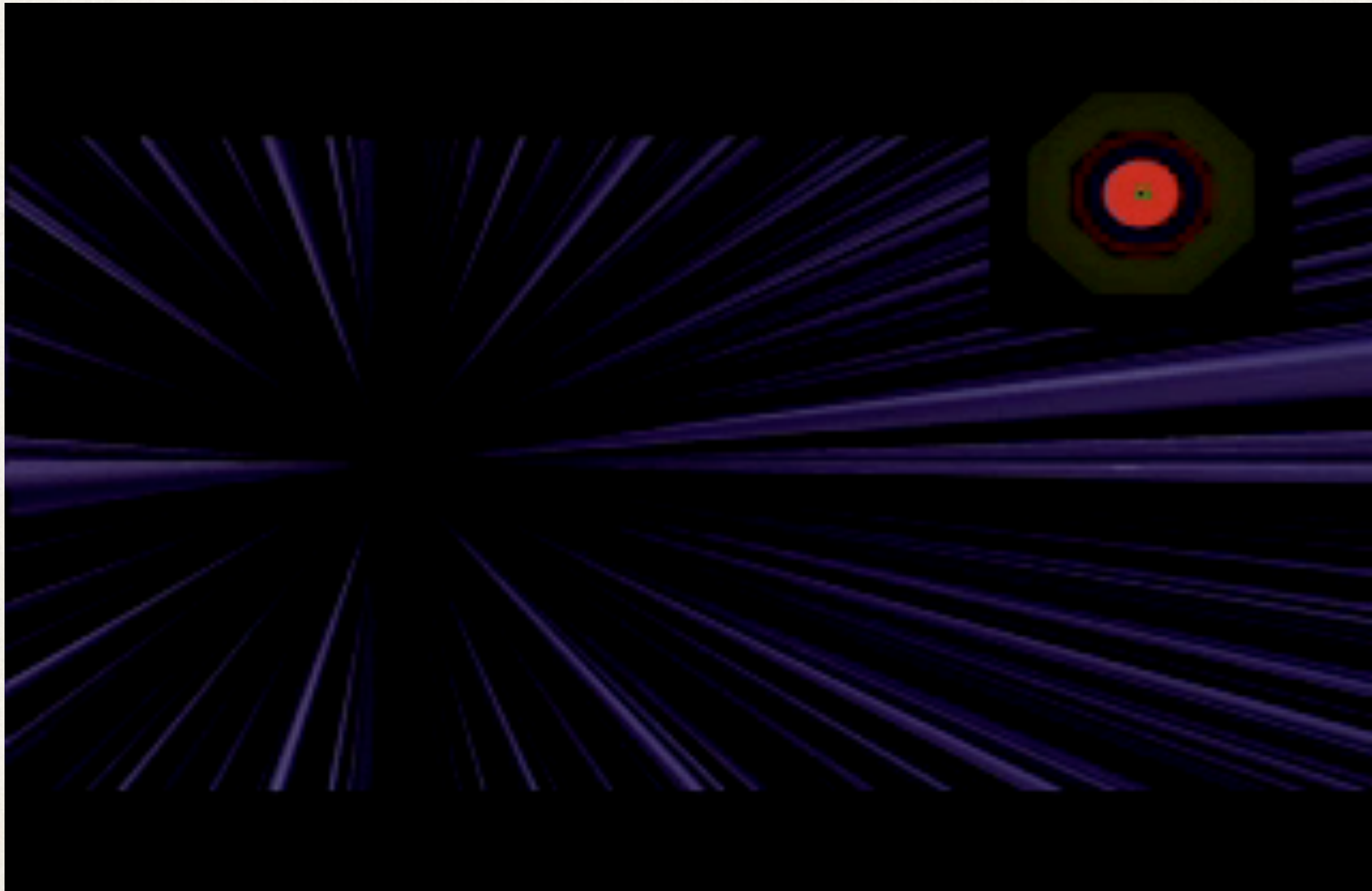
Detectores de gás

Schematic Principle of gas detectors



- Primary Ionization
- Secondary Ionization (due to δ -electrons)

Detectores de gás



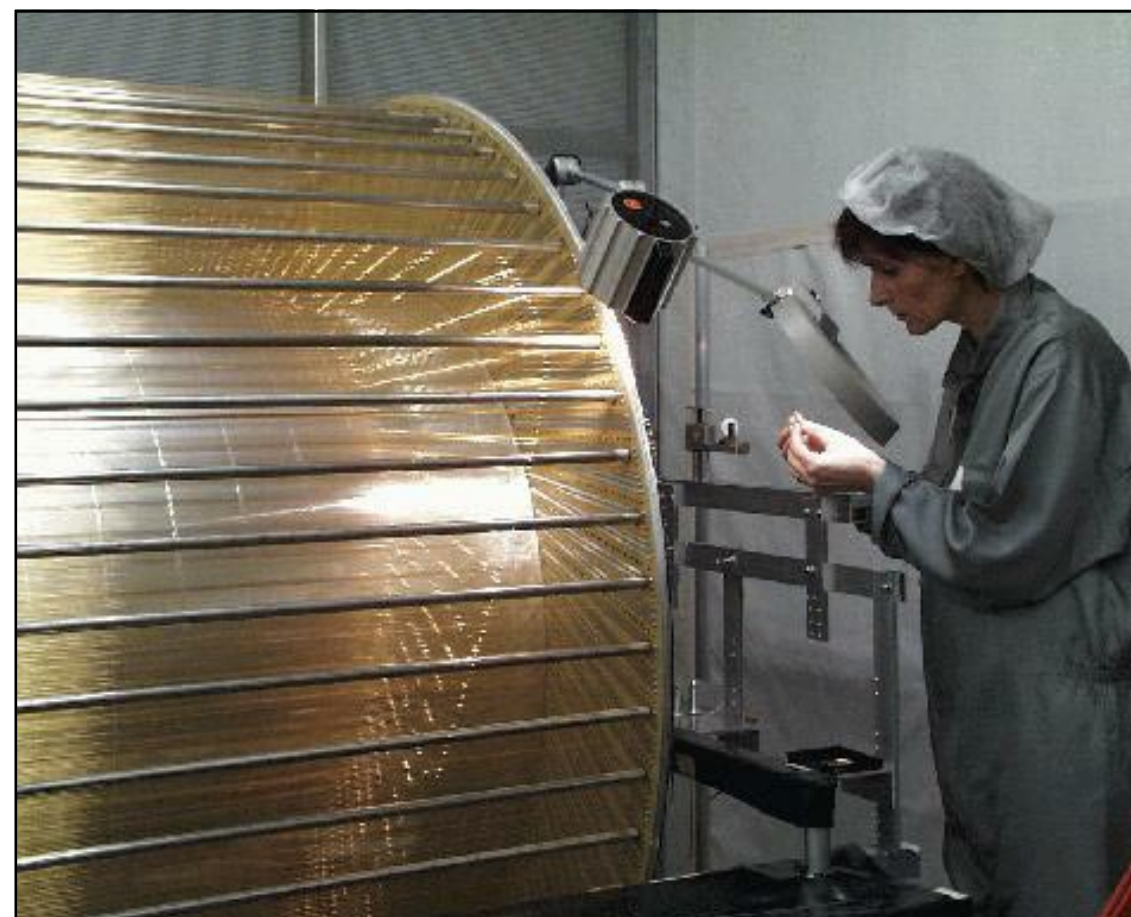
<http://cern.ch/amoraes>

Detectores de gás

6. Cylindrical Drift Chamber

Cylindrical
Drift Chamber

Number of wires: ~ 15000
Total force from wire tension: $\sim 6\text{ t}$



Detectores de gás

ALICE TPC (*Time Projection Chamber*)



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Detectores de gás

Relevant Parameters for gas detectors

Ionization energy

: E_i

Average energy/ion pair

: W_i

Average number of primary ion pairs [per cm]

: n_p

Average number of ion pairs [per cm]

: n_T

Differences
due to δ -electrons

$$\langle n_T \rangle = \frac{L \cdot \left\langle \frac{dE}{dx} \right\rangle_i}{W_i}$$

[about 2-6 times n_p]

[L: layer thickness]

δ -electrons lead to secondary ionization and limit spatial resolution; typical length scale of secondary ionization: 10 μm . Example: kinetic energy: $T_{\text{kin}} = 1 \text{ keV}$; gas: Isobutane $\rightarrow$ range: $R = 20 \mu\text{m} \dots$

[using $R [\text{g/cm}^2] = 0.71 (T_{\text{kin}})^{1.72} [\text{MeV}]$; valid for $T_{\text{kin}} < 100 \text{ keV}$]

Gas	$\langle Z \rangle$	$\rho [\text{g/cm}^3]$	$E_i [\text{eV}]$	$W_i [\text{eV}]$	$dE/dx [\text{keV/cm}]$	$n_p [\text{cm}^{-1}]$	$n_T [\text{cm}^{-1}]$
He	2	$1.66 \cdot 10^{-4}$	24.6	41	0.32	5.9	7.8
Ar	18	$1.66 \cdot 10^{-3}$	15.8	27	2.44	29.4	94
CH ₄	19	$6.7 \cdot 10^{-4}$	13.1	28	1.48	18	53
C ₄ H ₁₀	34	$2.42 \cdot 10^{-3}$	10.6	23	4.50	46	195

Detectores de gás

Ionization statistics:

Mean distance between two ionizations: $\lambda = 1/(n_e \sigma_I)$

Mean number of ionizations: $\langle n_p \rangle = L/\lambda$

n_p Poissonian distributed:

$$P(n_p, \langle n_p \rangle) = \frac{\langle n_p \rangle^{n_p} e^{-\langle n_p \rangle}}{n_p!}$$

σ_I : Ionization x-Section
 n_e : Electron density
 L : Thickness

Mean free path λ :
 [typical values]

He	0.25 cm
Air	0.052 cm
Xe	0.023 cm
[→ $\sigma_I(\text{He}) \approx 100 \text{ b}$]	

$P(0) = \exp(-L/\lambda)$ yields λ , σ_I
 using (in)efficiency of gas-detectors

Also important:

Mobility of charges:

Influences the timing behavior of gas detectors ...

Diffusion:

Influences the spatial resolution ...

Avalanche process via impact ionization:

Important for the gain factor of the gas detector ...

Recombination and electron attachment:

Admixture of electronegative gases (O_2 , F, Cl ...) influences detection efficiency ...

Deslocamento e difusão no gás

Ion mobility:

With external electric field: ions obtain velocity v_D in addition to thermal motion; on average ions move along field lines of electric field E ...

Kinetic energy:

$$\langle T_{\text{ion}}(E \neq 0) \rangle = \langle T_{\text{ion}}(\text{Therm.}) \rangle = \frac{3}{2}kT$$

Temperature
sorry ...

approximately equal to thermal energy, as the (heavy) ions loose typically half their energy when colliding with the non-ionized gas atoms.

Drift velocity v_D develops only from one interaction to another ...
Assuming $v_D(t=0)=0$ and collision time τ yields:

$$\vec{v} = \vec{a} \cdot \tau = \frac{e\vec{E}}{M} \cdot \tau$$

$$\vec{v}_D = \langle \vec{v} \rangle = \frac{1}{2}\vec{v} = \frac{e|\vec{E}|}{2M} \cdot \tau = \mu_+ |\vec{E}|$$

$$\tau = \lambda(T_{\text{kin}})/v_{\text{therm.}} = \text{const.}$$

since T_{kin} essentially thermal,
and $v_{\text{therm.}}$ thus constant ...

Drift velocity v_D for ions
proportional to E !

μ_+ : ion mobility e.g. $\mu_+=0.61 \text{ cm}^2/\text{Vs}$ for C_4H_{10}

[$E = 1 \text{ kV/cm}$; typical drift distances = few cm $\rightarrow$ typical ion drift time = few ms]

Deslocamento e difusão no gás

Electron mobility: $\vec{v}_D = \mu \vec{E}$
[B = 0]

Compare:

Electrons: v_D of order cm/ μ s

Ions: v_D of order cm/ms

Consider two situations:

$T_{\text{kin,e}} \gg kT$

gas atoms have only a few low-lying energy levels such that electrons can lose little energy in collisions [hot gases]

$$\lambda(T_e) \sim \lambda(E) \text{ and } \mu \sim \tau \sim 1/\sigma(E) \quad \begin{array}{l} \mu \text{ not constant!} \\ \text{[If } \lambda \sim 1/E; v_D = \text{const}] \end{array}$$

Electrons accelerated in E-field until sufficient energy is reached ...
Higher E-field yields smaller mean free path $\rightarrow$ constant v_D possible ...
[Example: $v_D = 3 - 5$ cm/ μ s for 90% Ar/10% CH₄]

$T_{\text{kin,e}} \approx kT$

gas atoms have many low-lying energy levels such that electrons lose all energy they gain between collisions [cold gases]

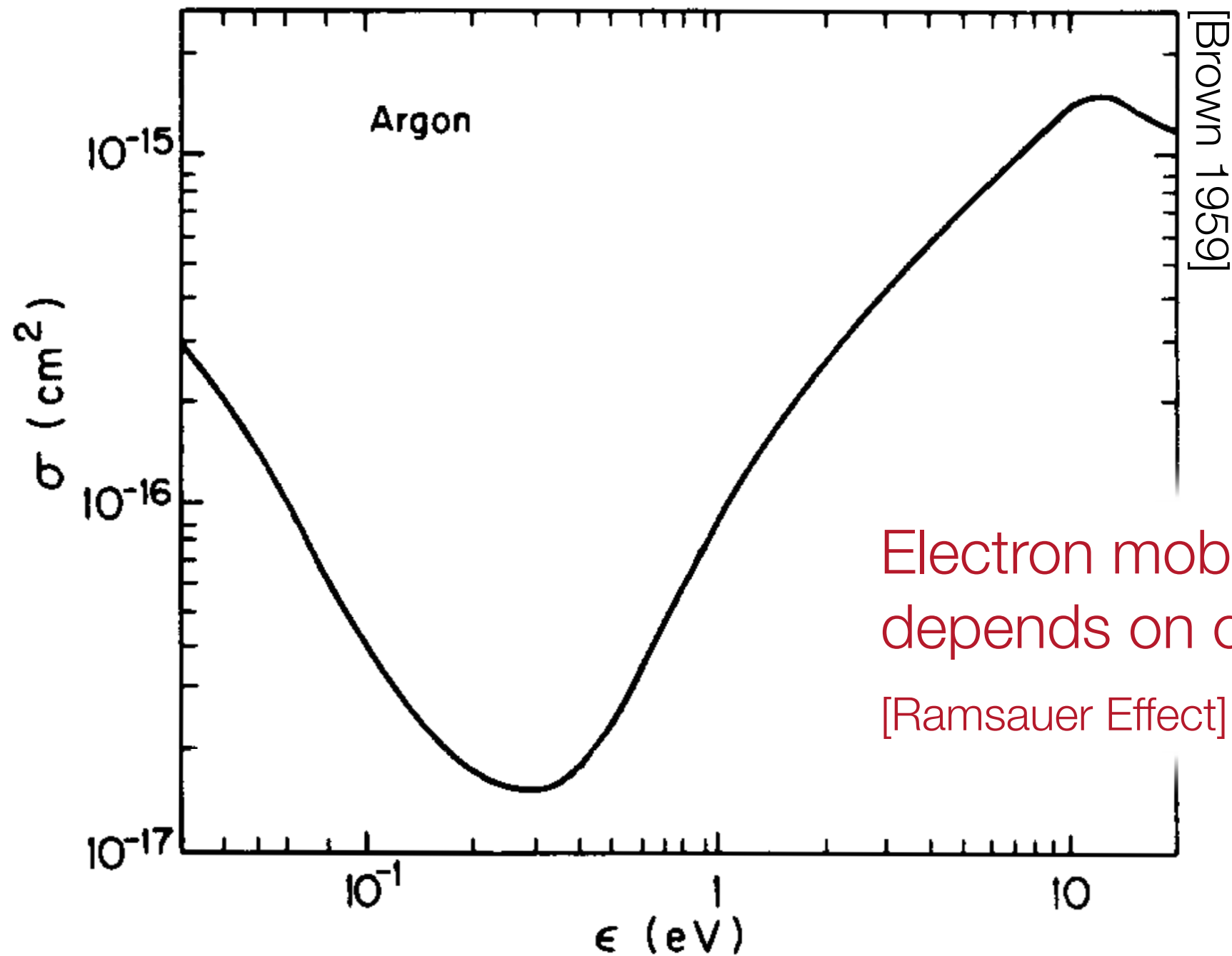
$$\mu \approx \text{const.} \quad \text{and} \quad v_D \propto E$$

Similar to situation with ions ...

[Example: $\mu = 7 \cdot 10^{-3}$ cm²/ μ s V for 90% Ne/10% CO₂; $v_D = 2$ cm/ μ s @ 300 V/cm]

Deslocamento e difusão no gás

Ramsauer Cross Section [from Sauli 1977]



Electron mobility
depends on cross section
[Ramsauer Effect]

Perda de elétrons

Electrons maybe lost during drift ...

Possible processes:

i. recombination of ions and electrons

Depends on number of charge carriers
and recombination coefficient ...

Generally not important ...

Recombination rate:

$$\Lambda = p_r \cdot n^+ n^-$$

➡ Recombination
coefficient $\approx 10^{-7} \text{ cm}^3/\text{s}$

ii. electron attachment

Electro-negative gases bind electrons; e.g.: O_2 , Freon, Cl_2 , SF_6 ...

Attachment coefficient h strongly energy dependent ("Ramsauer effect") ...

Example O_2 : $h = 10^{-4}$

Collisions of electron per second: 10^{11}

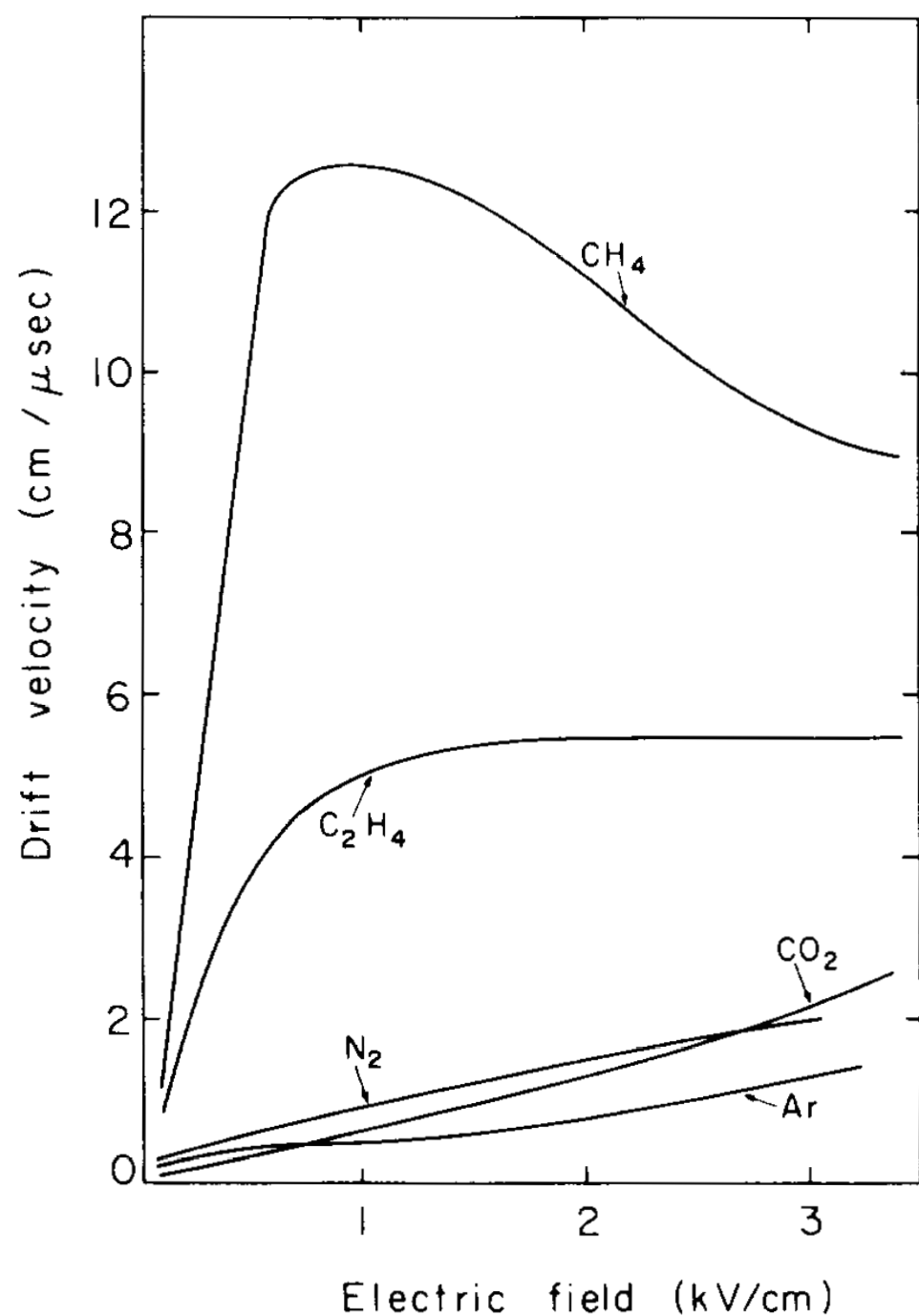
Typical drift time of electron: 10^{-6} s

Fraction lost: $X_{\text{loss}} = 10^{-4} 10^{11} \text{ s}^{-1} 10^{-6} \text{ s} \cdot p = 10p$

$X_{\text{loss}} < 1\% \rightarrow p < 10^{-3}$, i.e. less than 1 ‰ admixture

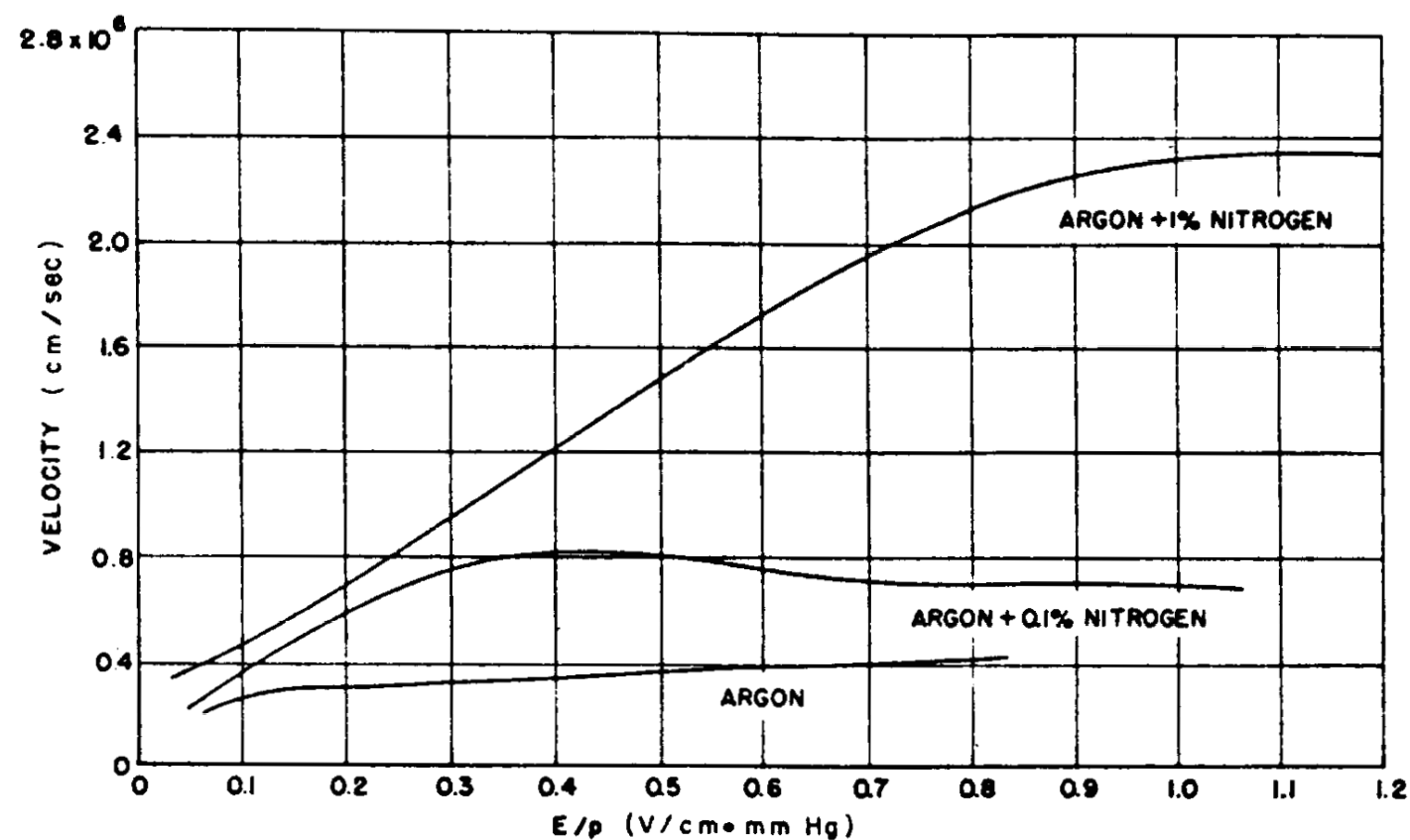
Oxygen should
be kept out

Deslocamento e difusão no gás



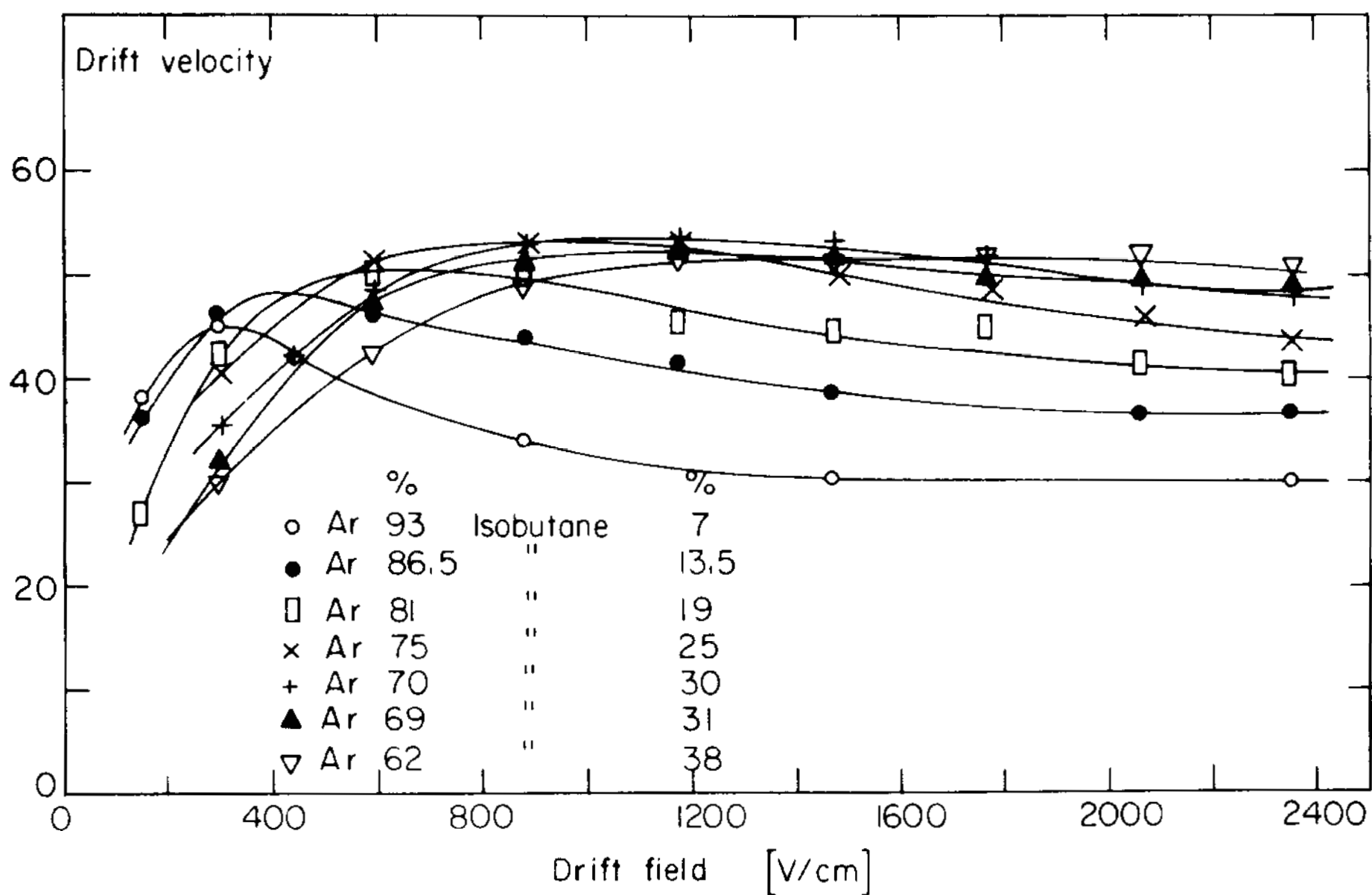
Drift velocity of electrons
in several gases at normal conditions

Use gas mixture to obtain constant v_D
Important for applications using drift time to get
spatial information



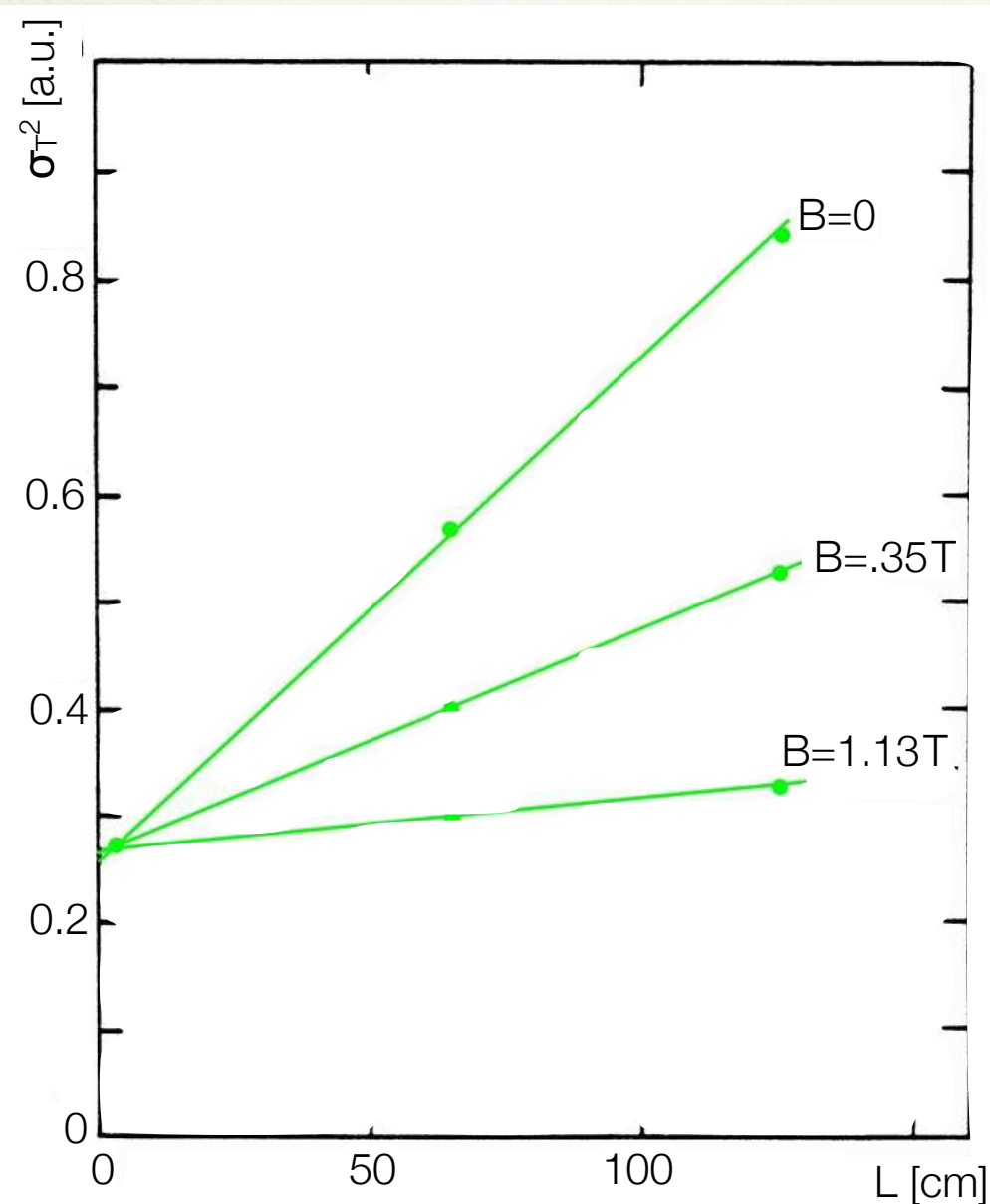
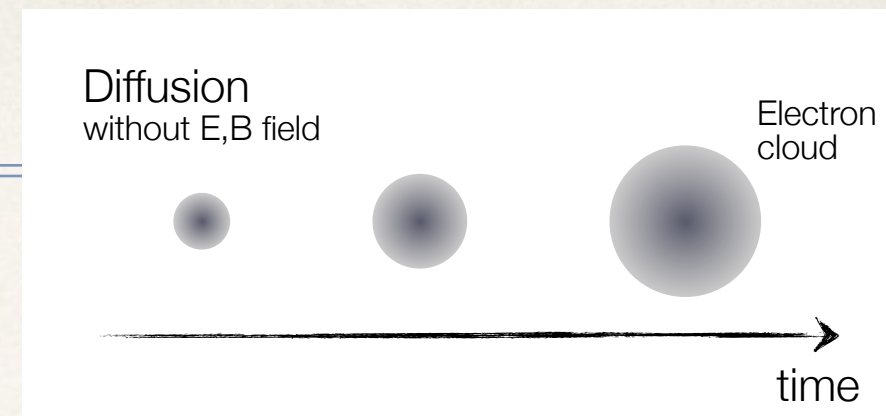
E-Field/pressure

Deslocamento e difusão no gás

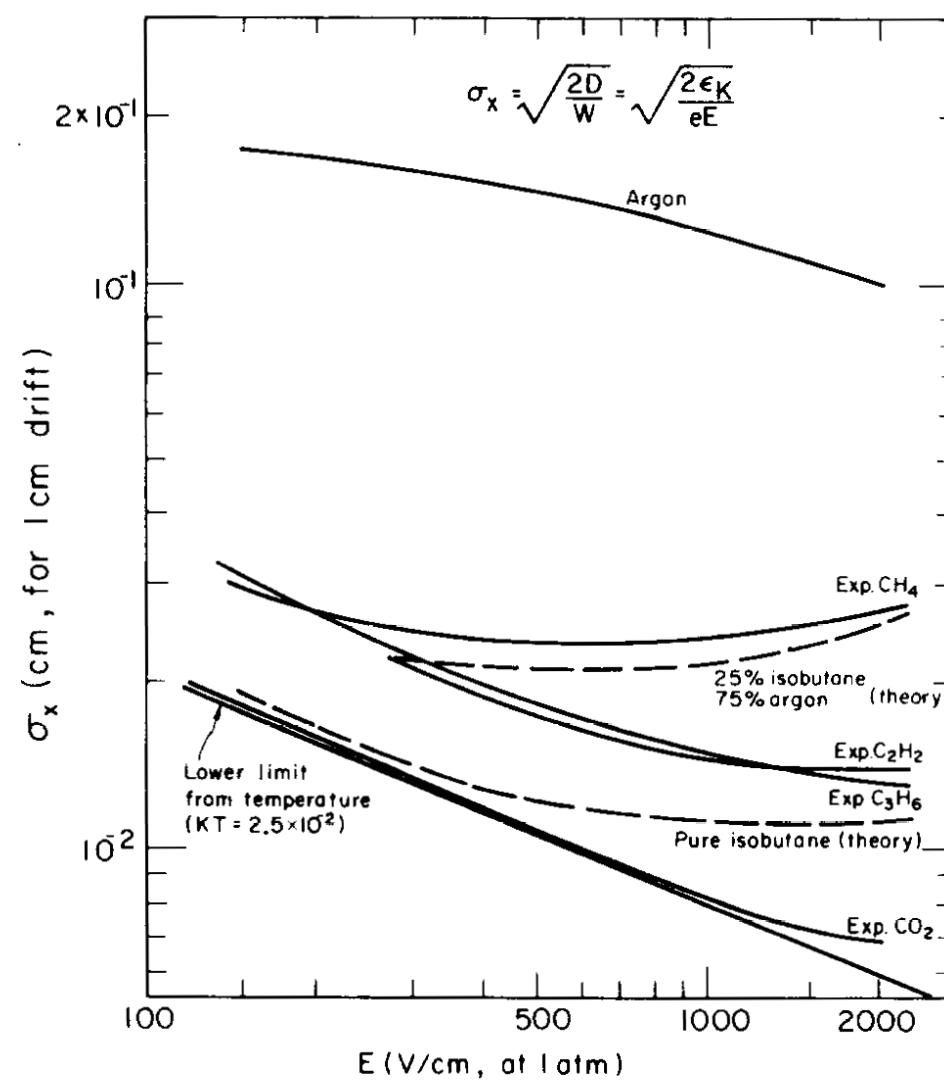


Drift velocity in several argon-isobutane (C₄H₁₀) mixtures

Deslocamento e difusão no gás



Longitudinal diffusion in E-field



Multiplicação em avalanche

Large electric field yields
large kinetic energy of electrons ...

→ Avalanche formation

Larger mobility of electrons results in liquid
drop like avalanche with electrons near head ...

Mean free path: λ_{ion}
[for a secondary ionization]

Probability of an ionization per
unit path length: $\alpha = 1/\lambda_{\text{ion}}$ [1st Townsend coefficient]

$$dn = n \cdot \alpha dx$$

$n(x)$ = electrons
at location x

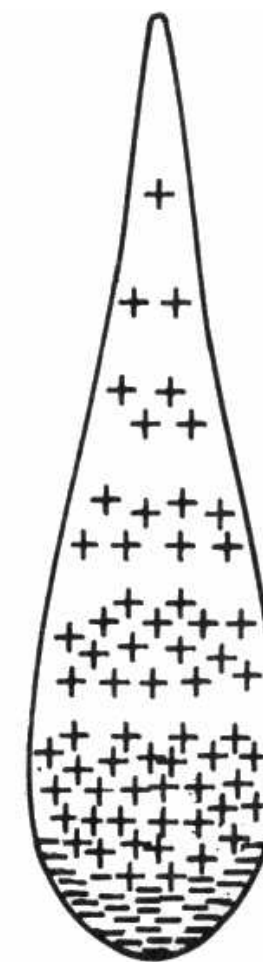
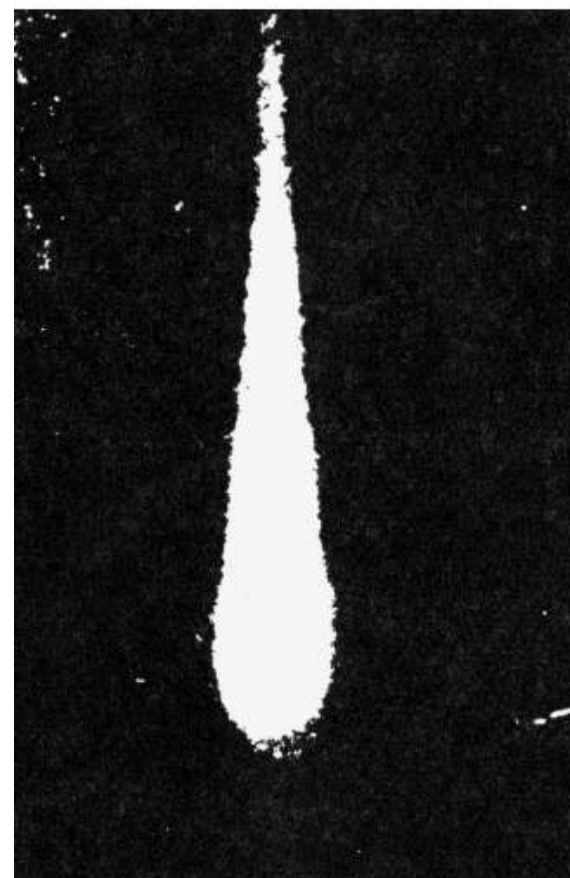
$$n = n_0 e^{\alpha x}$$

Gain:

$$G = \frac{n}{n_0} = e^{\alpha x} \quad \text{and more general for } \alpha = \alpha(x): \quad G = \frac{n}{n_0} = \exp \left[\int_{x_1}^{x_2} \alpha(x) dx \right]$$

[Raether limit: $G \approx 10^8$; $\alpha x = 20$; then sparking sets in ...]

Townsend avalanche



Drop-like shape of an avalanche

Left: cloud chamber picture

Right: schematic view

Multiplicação em avalanche

To reach high E fields
use a thin wire as anode ...

Close to wire E-field very large ...

Kinetic energy:

$$\Delta T_{\text{kin}} = e\Delta U = e \int_{r_1}^{r_2} E(r) dr$$

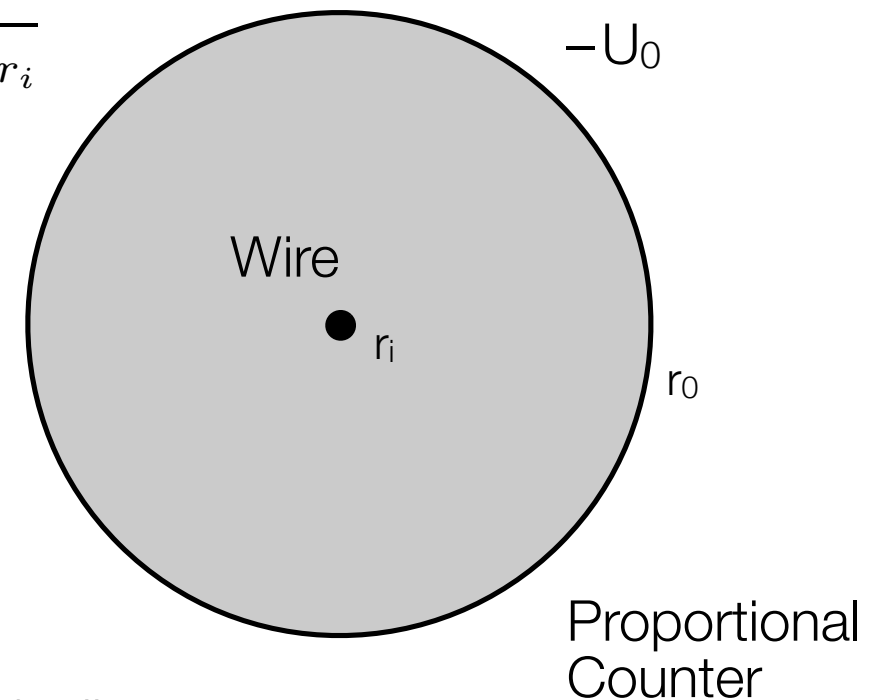
$$= \frac{eU_0}{\ln r_0/r_i} \int_{r_1}^{r_2} \frac{1}{r} dr$$

$$= \underbrace{\frac{eU_0}{\ln r_0/r_i}}_{\text{const.}} \cdot \ln r_2/r_1$$

Choose ratio as
large as possible with small Δr ...

[Δr should be smaller than typical mean free path ...]

$$E = \frac{U_0}{r \ln r_0/r_i}$$



Reminder Physics II:

$$\text{div} \vec{E} = \frac{\rho}{\epsilon_0}$$

$$\rightarrow |\vec{E}| = \frac{\lambda}{2\pi\epsilon_0} \frac{1}{r}$$

$$\oint \vec{E} dA = \int \frac{\rho}{\epsilon_0} d^3r$$

$$2\pi r L \cdot |\vec{E}| = Q/\epsilon_0$$

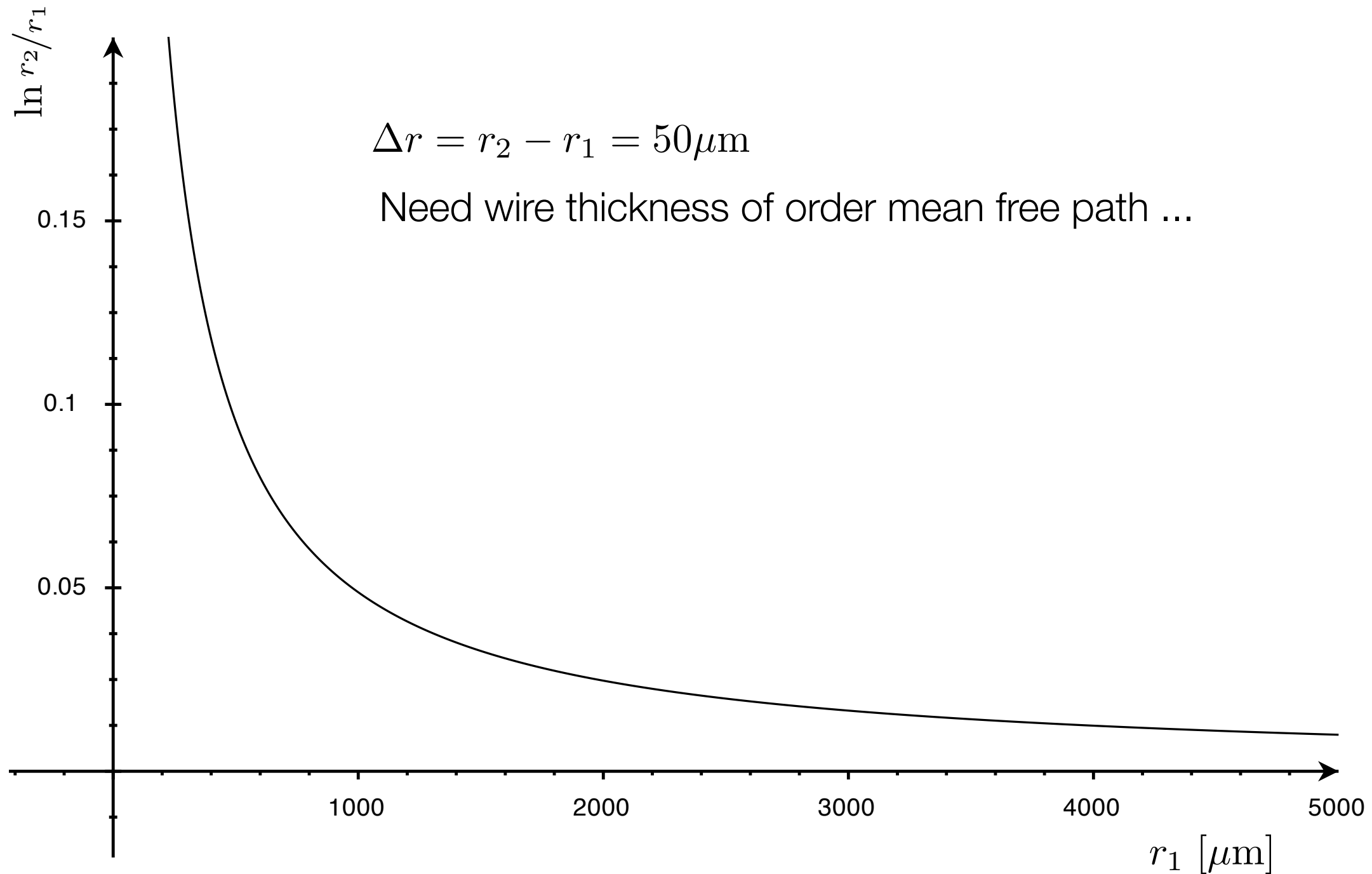
$$\lambda = Q/L$$

linear charge
density

$$U_0 = \int_{r_i}^{r_0} E dr = \frac{\lambda}{2\pi\epsilon_0} \ln \frac{r_0}{r_i}$$

$$\rightarrow \frac{\lambda}{2\pi\epsilon_0} = \frac{U_0}{\ln r_0/r_i}$$

Multiplicação em avalanche



Single wire proportional counter

Ionization mode:

full charge collection
no multiplication; gain ≈ 1

Proportional mode:

multiplication of ionization
signal proportional to ionization
measurement of dE/dx
secondary avalanches need quenching;
gain $\approx 10^4 - 10^5$

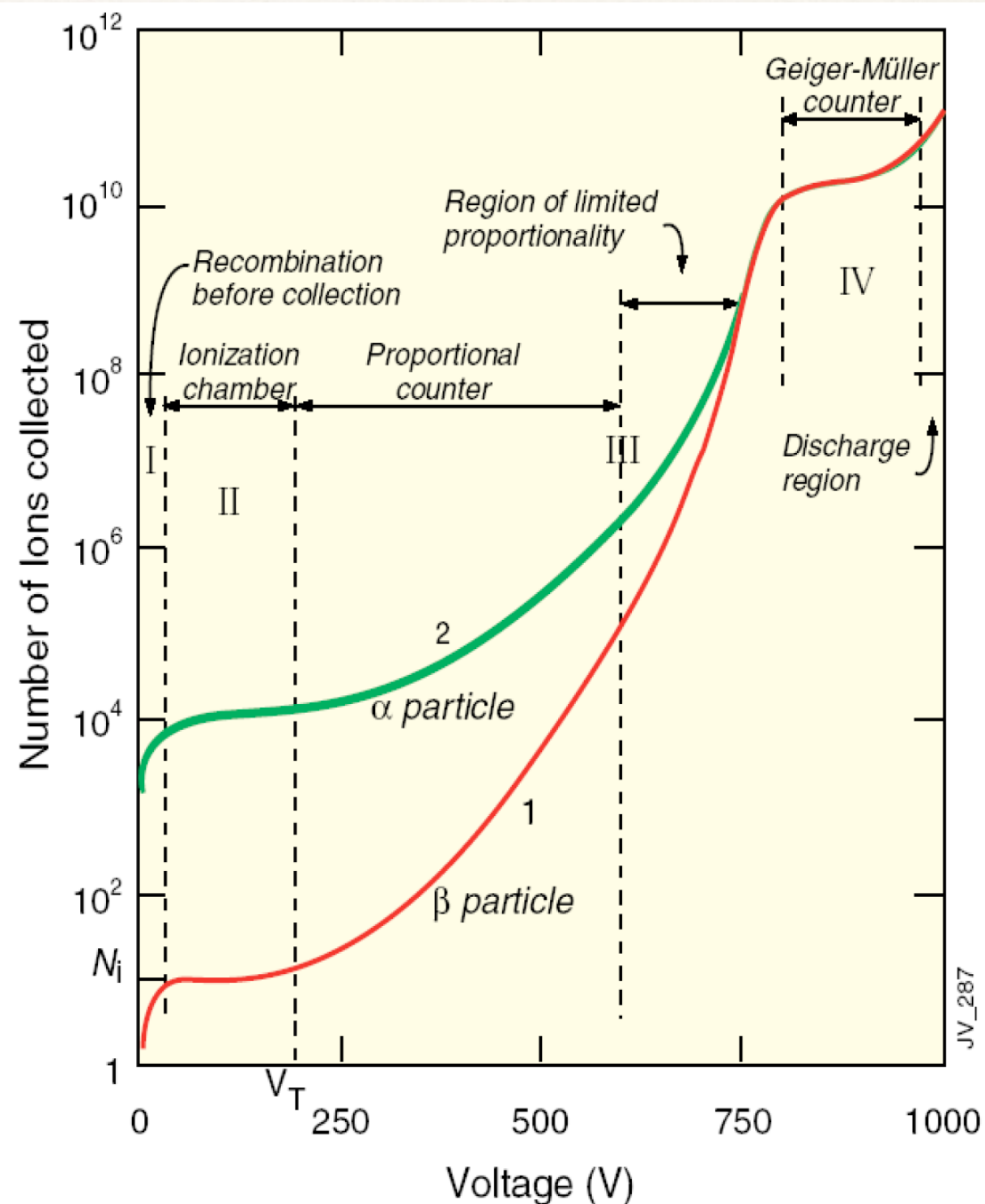
Limited proportional mode:

[saturated, streamer]

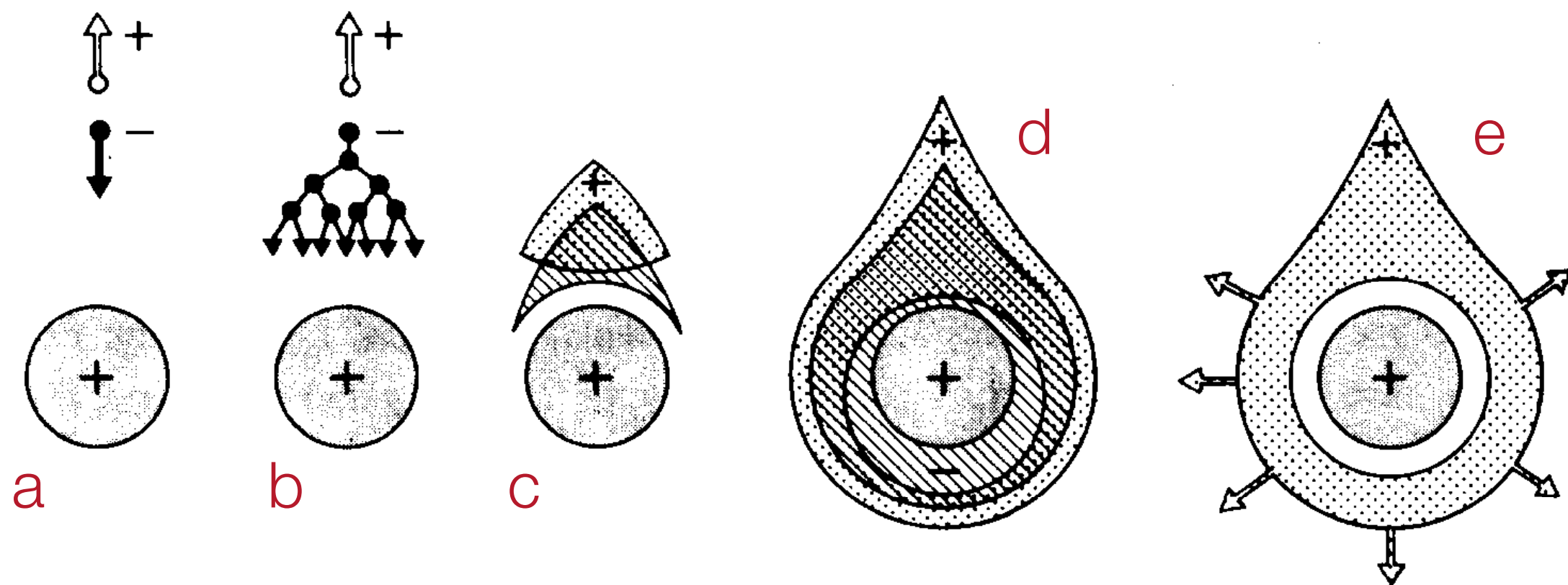
strong photoemission
requires strong quenchers or pulsed HV;
gain $\approx 10^{10}$

Geiger mode:

massive photoemission;
full length of the anode wire affected;
discharge stopped by HV cut



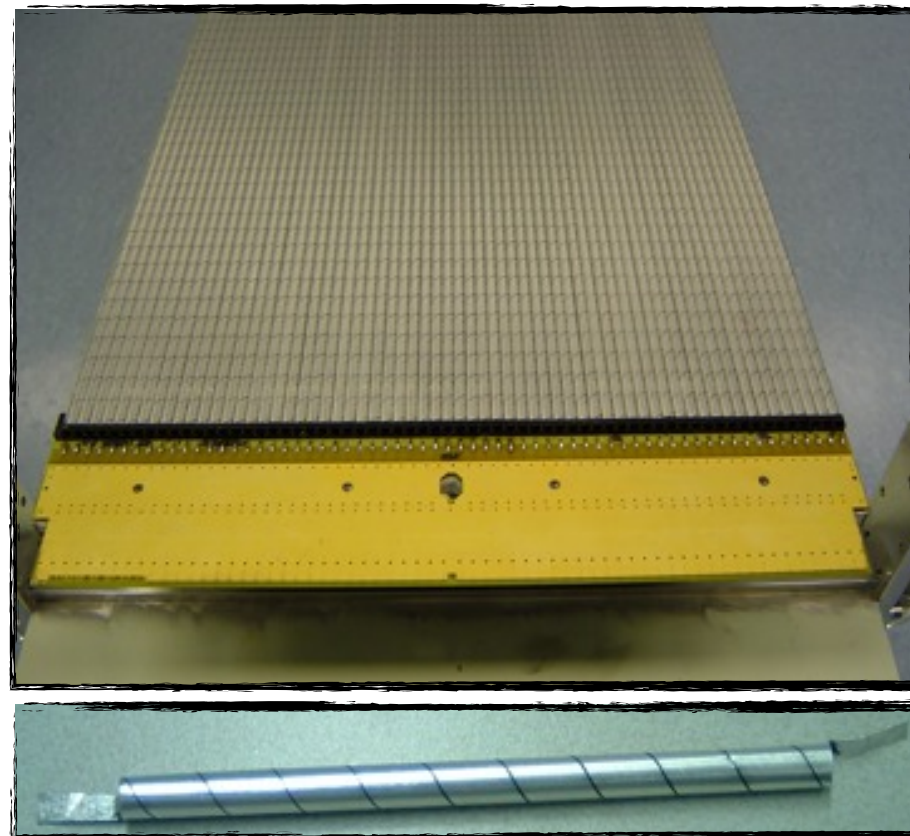
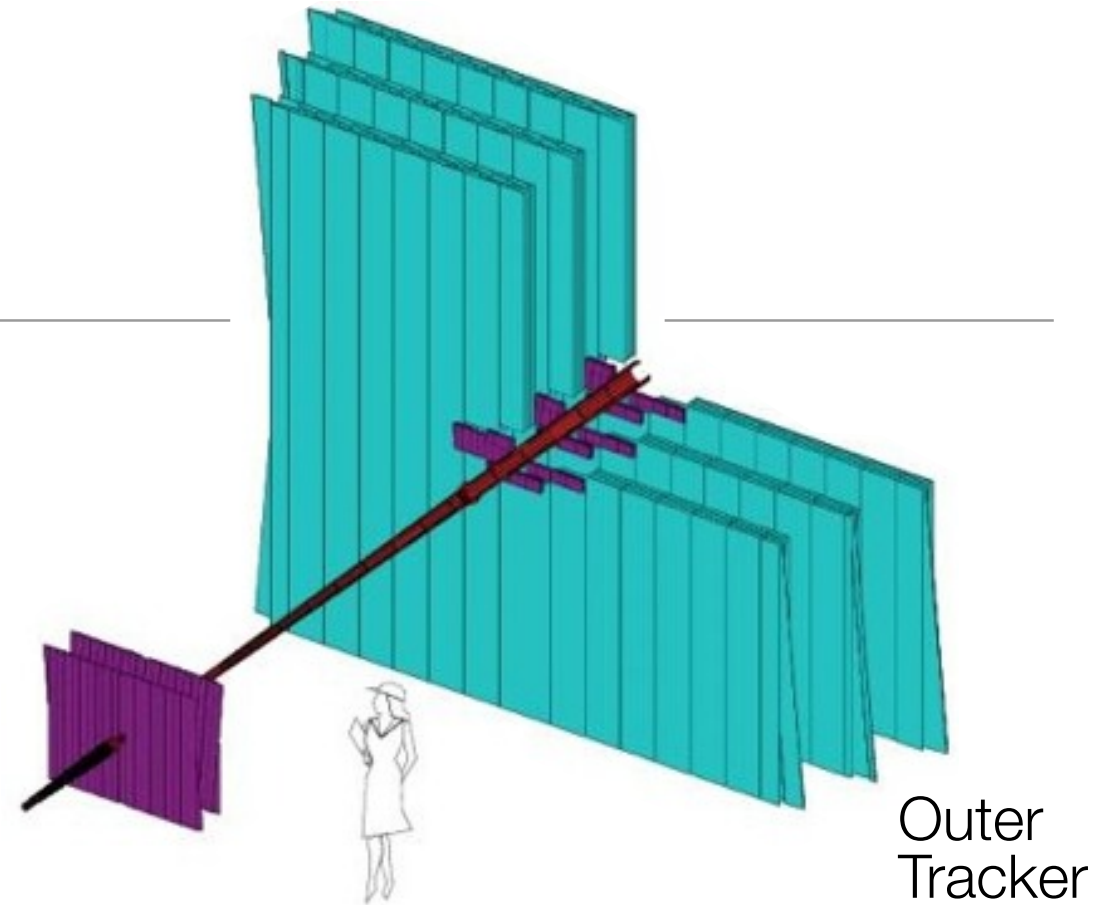
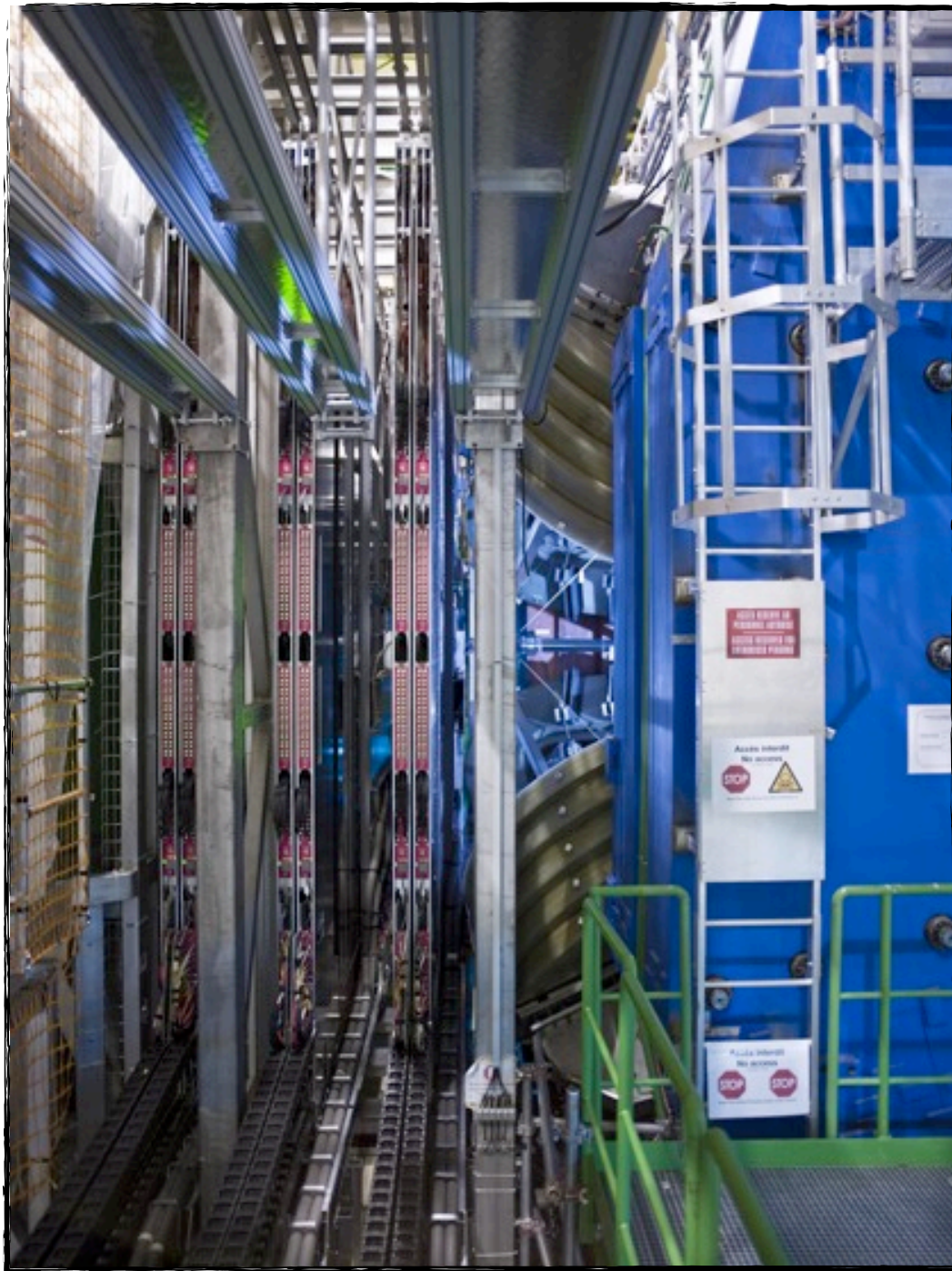
Deslocamento e difusão no gás



Time development of an avalanche in a proportional counter

A single primary electron proceeds towards anode in regions of increasingly high fields, experiencing ionizing collisions; due to the lateral diffusion, a drop-like avalanche, surrounding the wire develops.

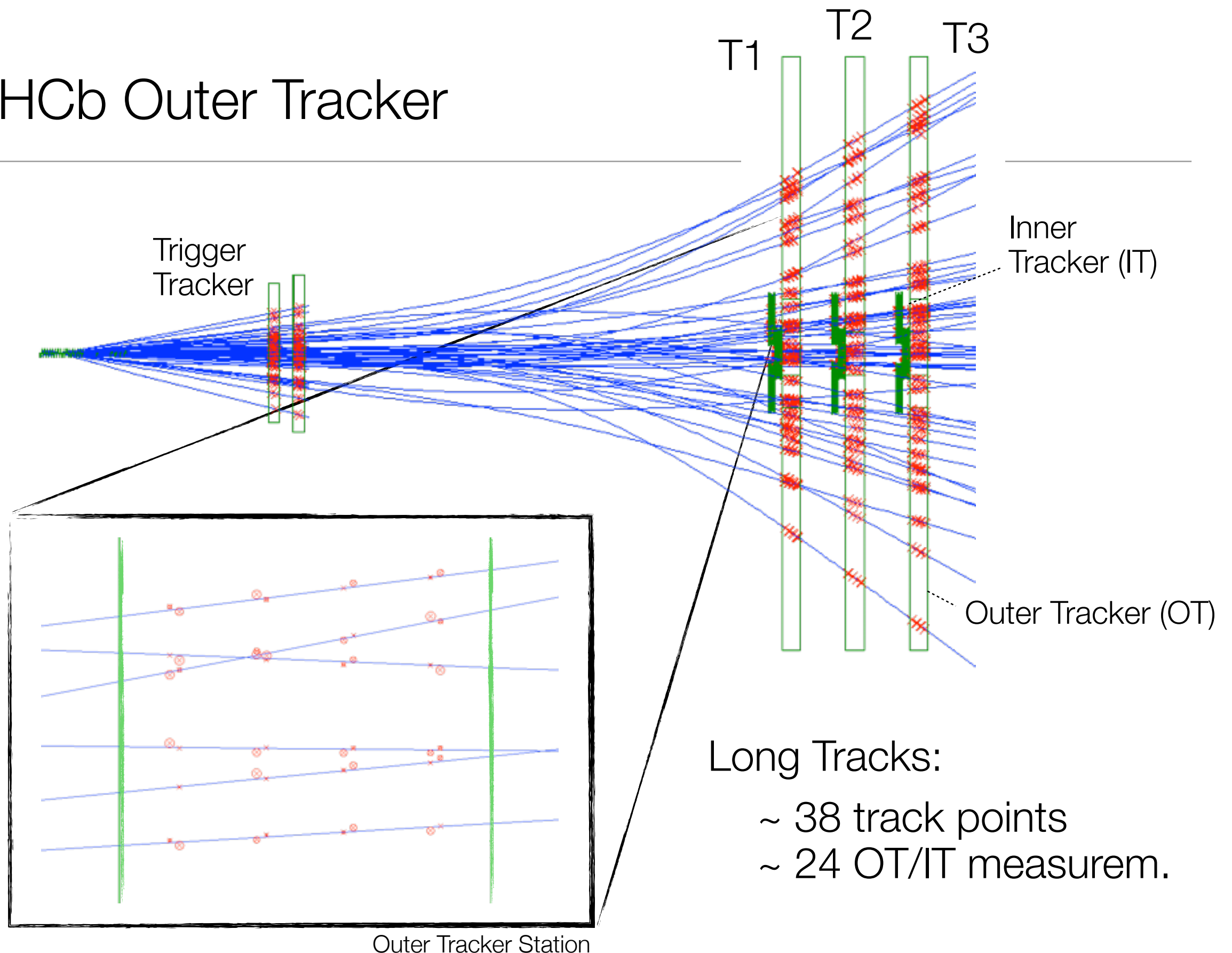
LHCb Outer Tracker



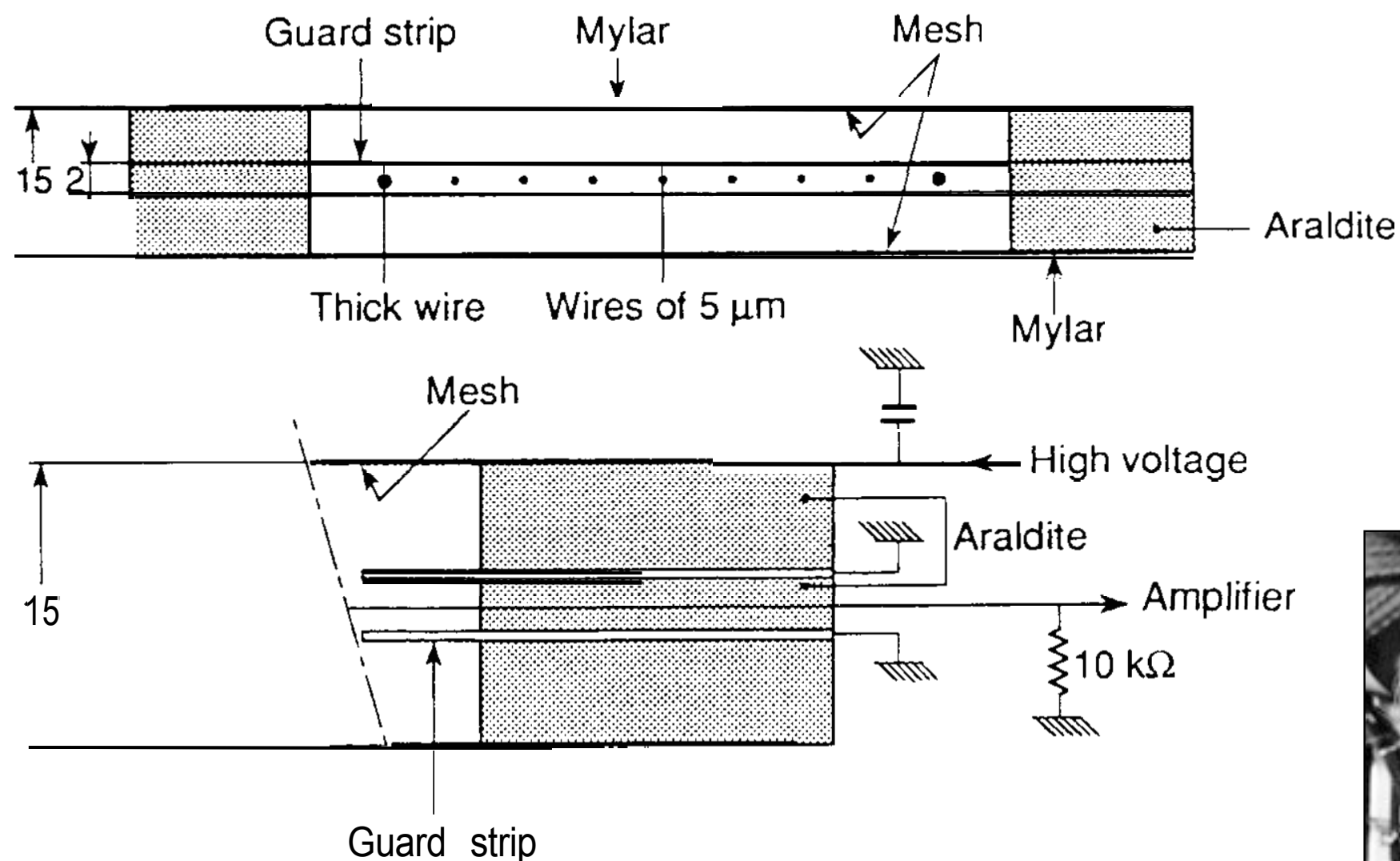
Straw Tubes
[double layers]

3 Chambers
[4 layers á 18 modules]

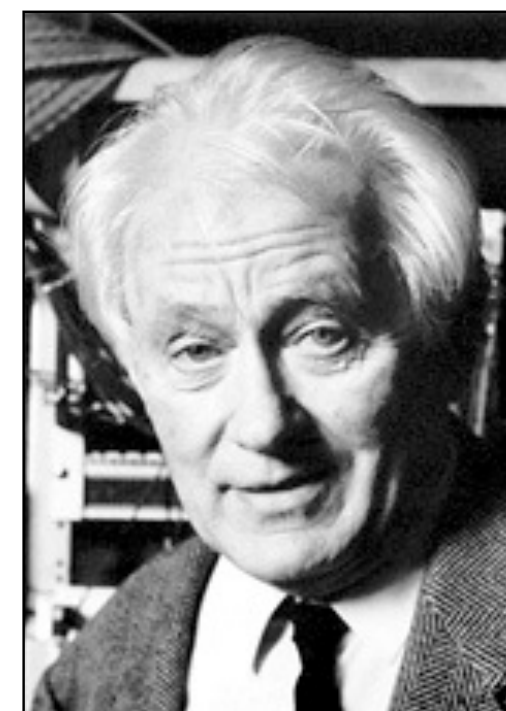
LHCb Outer Tracker



Multi-wire proportional chamber (MWPC)



G. Charpak
Nobel Prize 1992

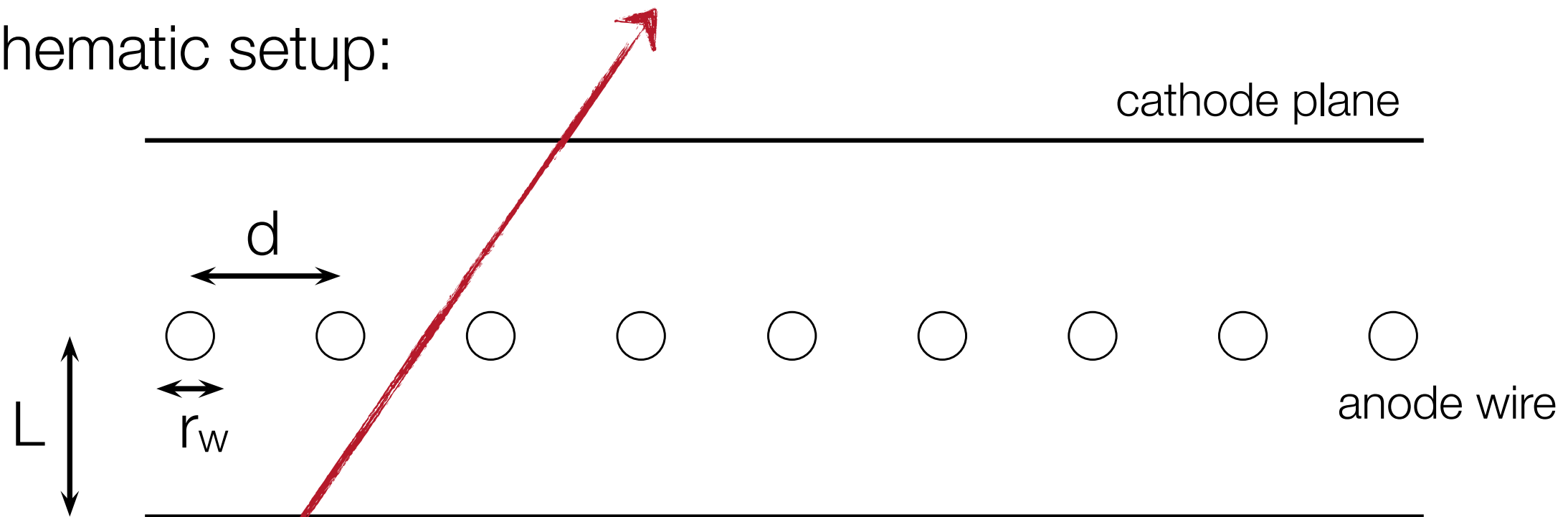


MWPC construction details
from Charpak's nobel lecture [1967 design]

Sense wires [$\varnothing = 20 \mu\text{m}$] separated by 2 mm; wires lie between two cathode meshes; edges of the planes are potted in Araldite ...

Multi-wire proportional chamber (MWPC)

Schematic setup:



Parameters:

$d = 2 - 4 \text{ mm}$
 $r_w = 20 - 25 \text{ } \mu\text{m}$
 $L = 3 - 6 \text{ mm}$
 $U_0 = \text{several kV}$
 Total area: $O(\text{m}^2)$

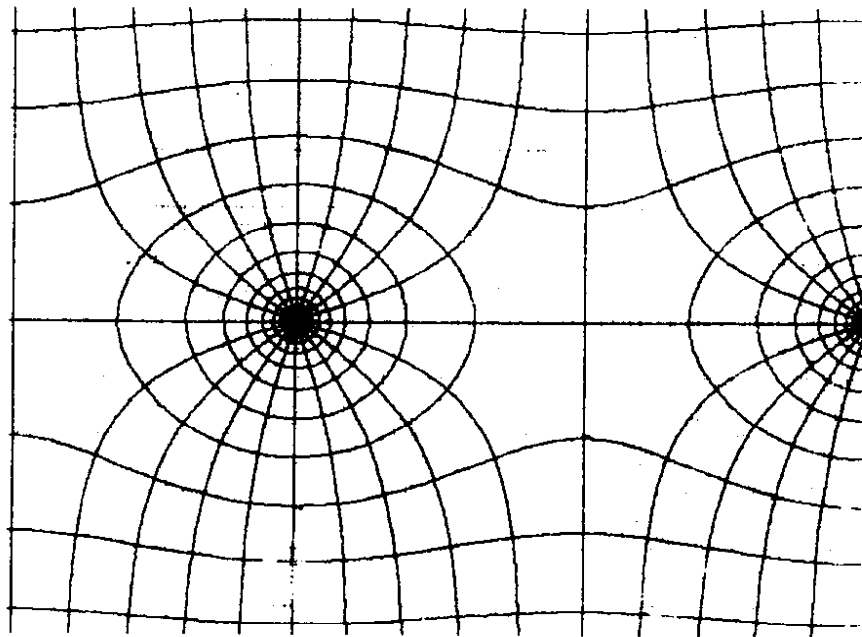
Features:

Tracking of charged particles
 Some PID capabilities via dE/dx
 Large area coverage
 High rate capabilities

particle track

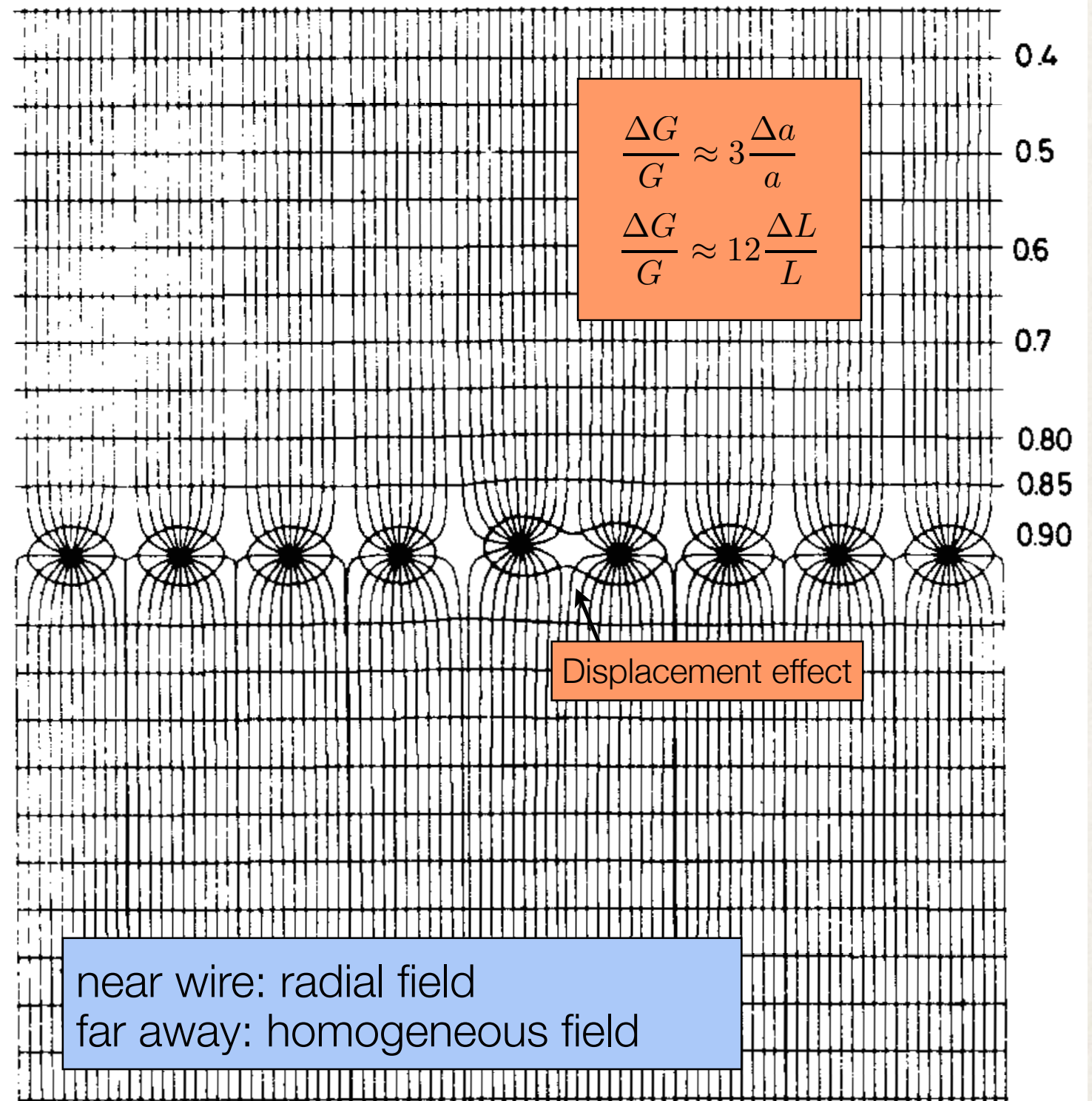
Multi-wire proportional chamber (MWPC)

Electric field lines
and equipotentials



Small wire displacements
reduce field quality ...

Need high mechanical precision
both for geometry and wire tension ...
[electrostatics and gravitation; wire sag]



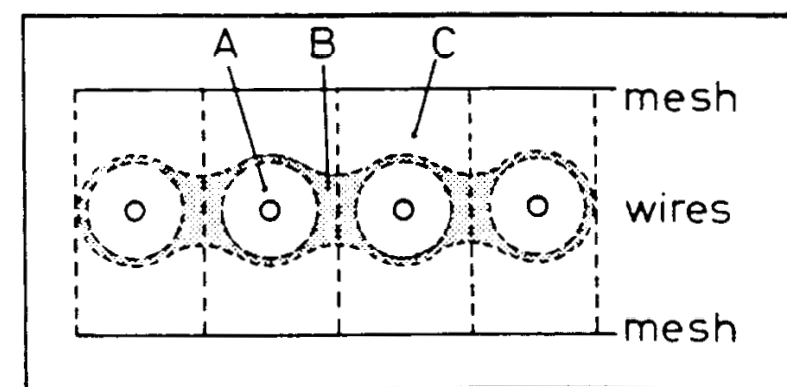
Multi-wire proportional chamber (MWPC)

Signal generation:

Electrons drift to closest wire

Gas amplification near wire $\rightarrow$ avalanche

Signal generation due to electrons and slow ions ...



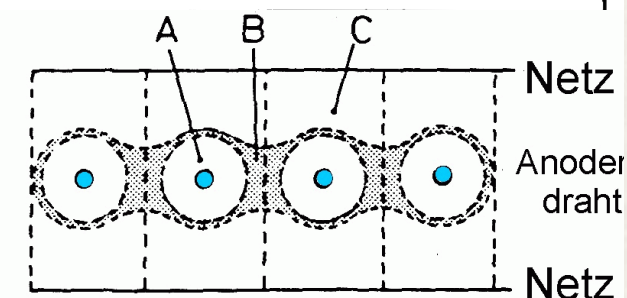
Timing resolution:

Depends on location of penetration

For fast response: OR of all channels .

[Typical: $\sigma_t = 10$ ns]

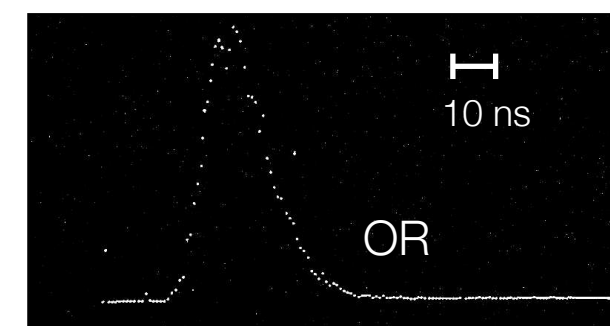
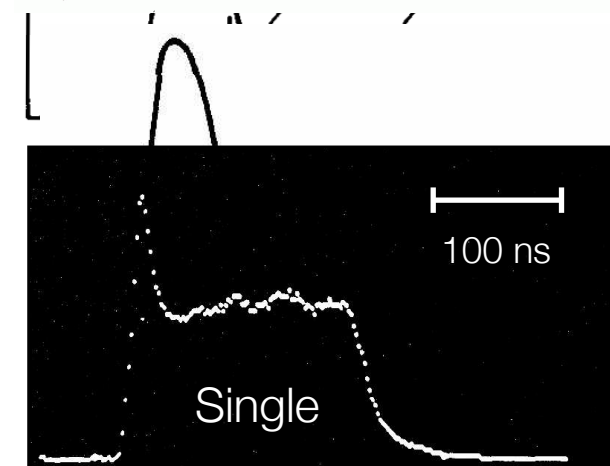
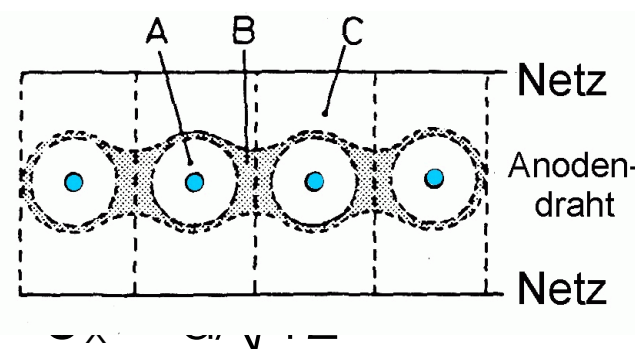
main
contribution



Space point resolution:

Only information about closest wire $\rightarrow$

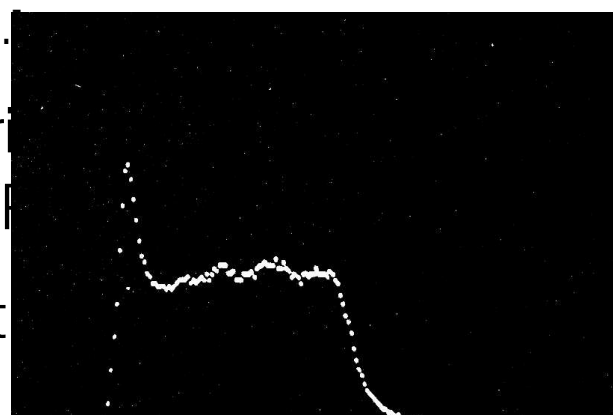
[Not very precise and only one for one dimension ...]



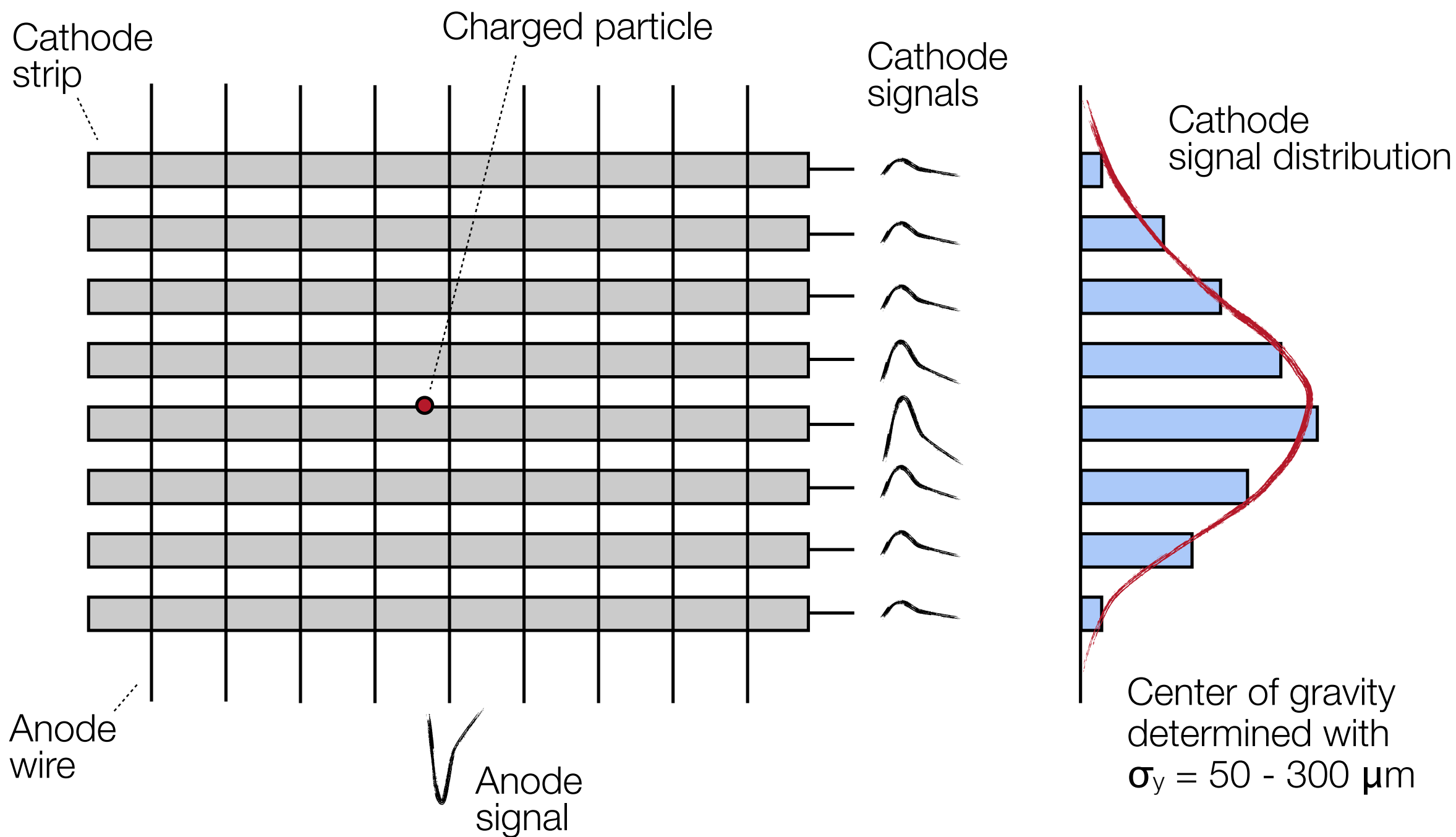
2-dim.: use 2 MWPCs with different orientation

3-dim.: several layers of such X-Y-MWPC

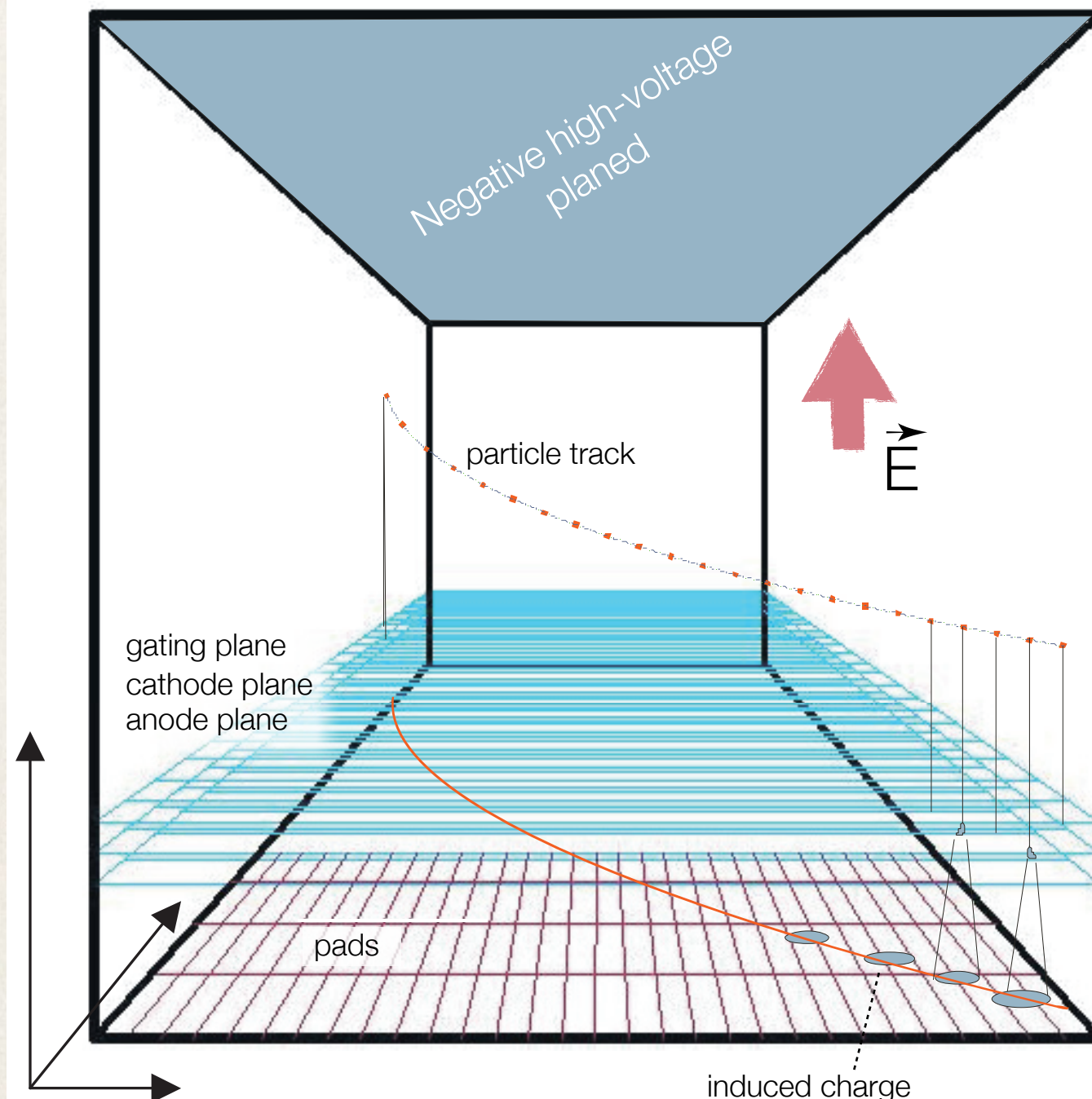
Possible improvement: segmented cathode



Multi-wire proportional chamber (MWPC)



Time projection chamber (TPC)



Advantages:

- Complete track within one detector yields good momentum resolution
- Relative few, short wires (MWPC only)
- Good particle ID via dE/dx
- Drift parallel to B suppresses transverse diffusion by factors 10 to 100

Challenges:

- Long drift time; limited rate capability [attachment, diffusion ...]
- Large volume [precision]
- Large voltages [discharges]
- Large data volume ...
- Extreme load at high luminosity; gating grid opened for triggered events only ...

Typical resolution:

z : mm; x : 150 - 300 μm ; y : mm
 dE/dx : 5 - 10%

Time projection chamber (TPC)

ALICE TPC:

Length: 5 meter

Radius: 2.5 meter

Gas volume: 88 m³

Total drift time: 92 μ s

High voltage: 100 kV

End-cap detectors: 32 m²

Readout pads: 557568

159 samples radially

1000 samples in time

Gas: Ne/CO₂/N₂ (90-10-5)

Low diffusion (cold gas)

Gain: $> 10^4$

Diffusion: $\sigma_t = 250 \mu\text{m}$

Resolution: $\sigma \approx 0.2 \text{ mm}$

$\sigma_p/p \sim 1\% \text{ p}; \epsilon \sim 97\%$

$\sigma_{dE/dx}/(dE/dx) \sim 6\%$

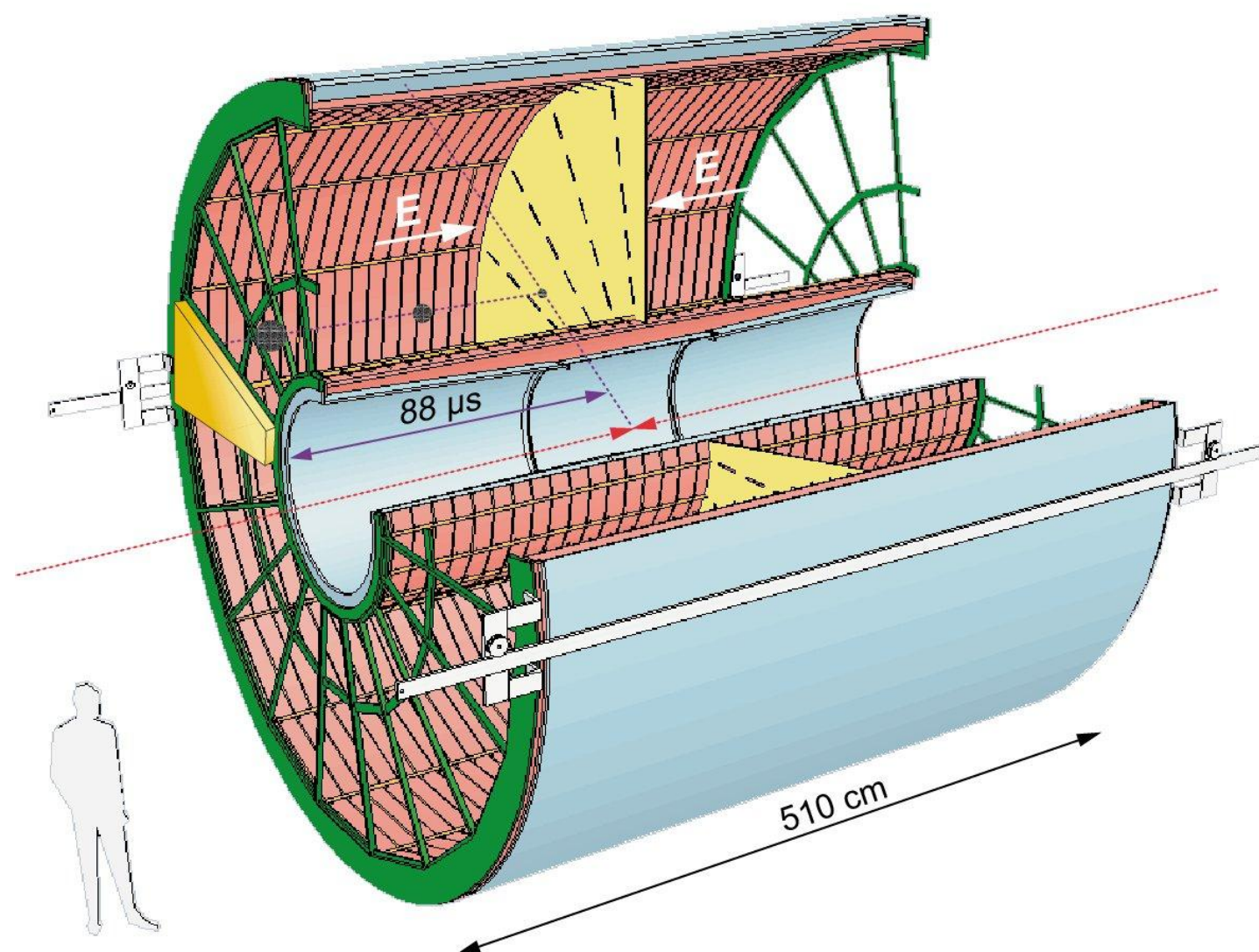
Magnetic field: 0.5 T

Pad size: 5x7.5 mm² (inner)

6x15 mm² (outer)

Temperature control: 0.1 K

[also resistors ...]



Material: Cylinder build from composite material of airline industry ($X_0 \sim 3\%$)

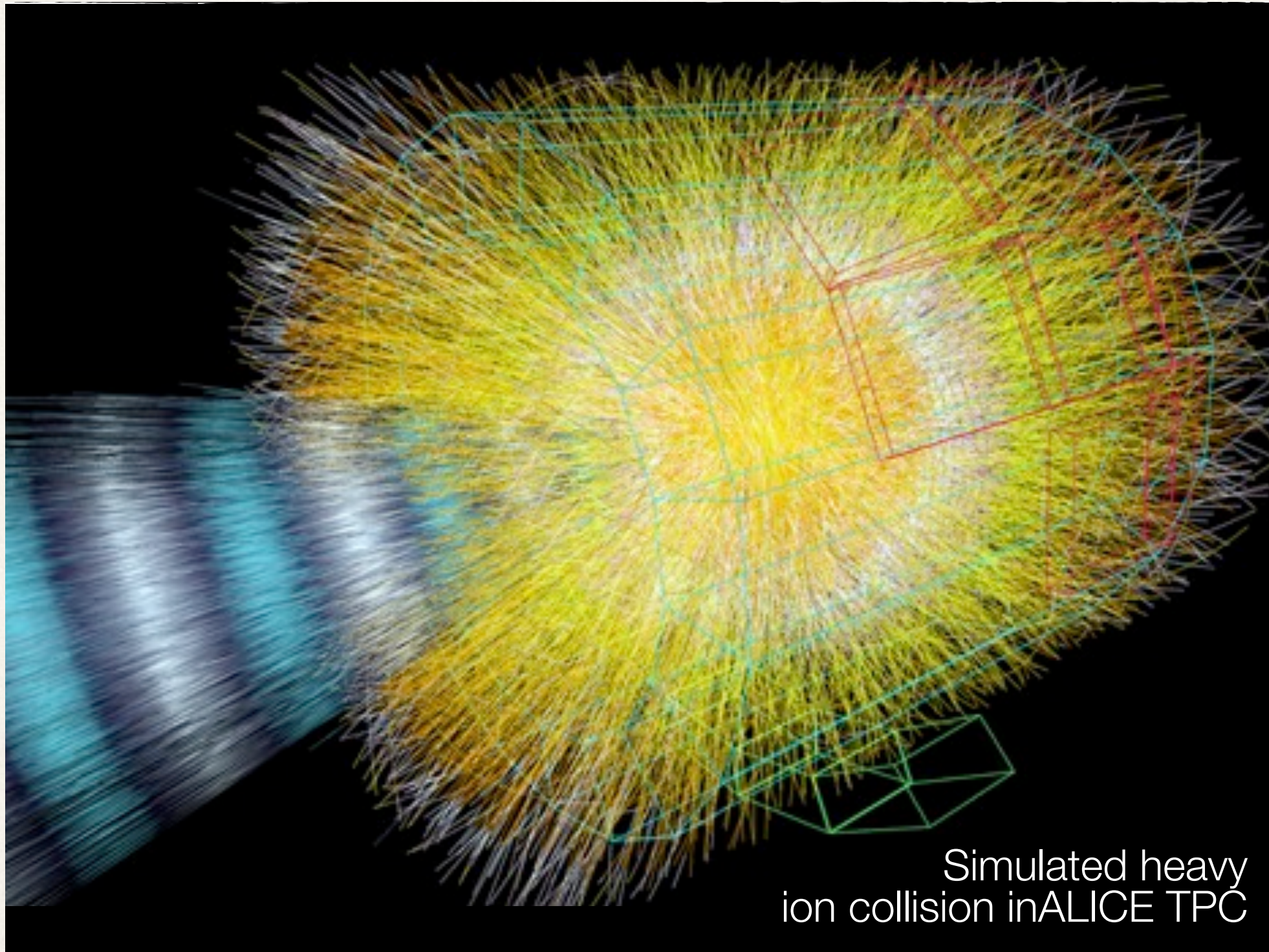
Time projection chamber (TPC)



View inside
ALICE TPC

<http://cern.ch/amoraes>

Time projection chamber (TPC)



Simulated heavy
ion collision in ALICE TPC

Envelhecimento das câmaras de fios

Avalanche formation can be considered as micro plasma discharge.

Consequences:

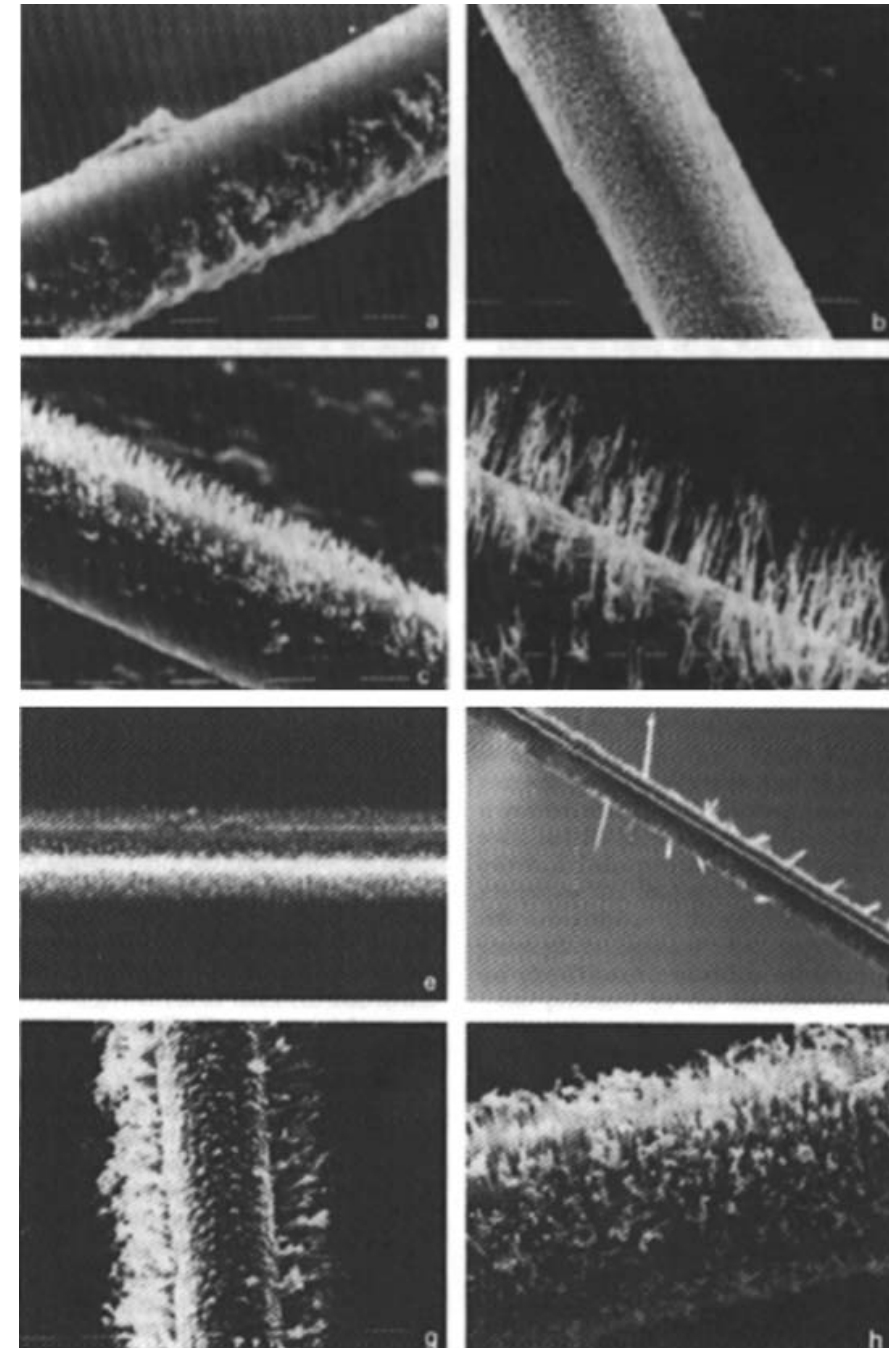
- Formation of radicals i.e. molecule fragments
- Polymerization yields long chains of molecules
- Polymers may be attached to the electrodes
- Reduction of gas amplification

Important:

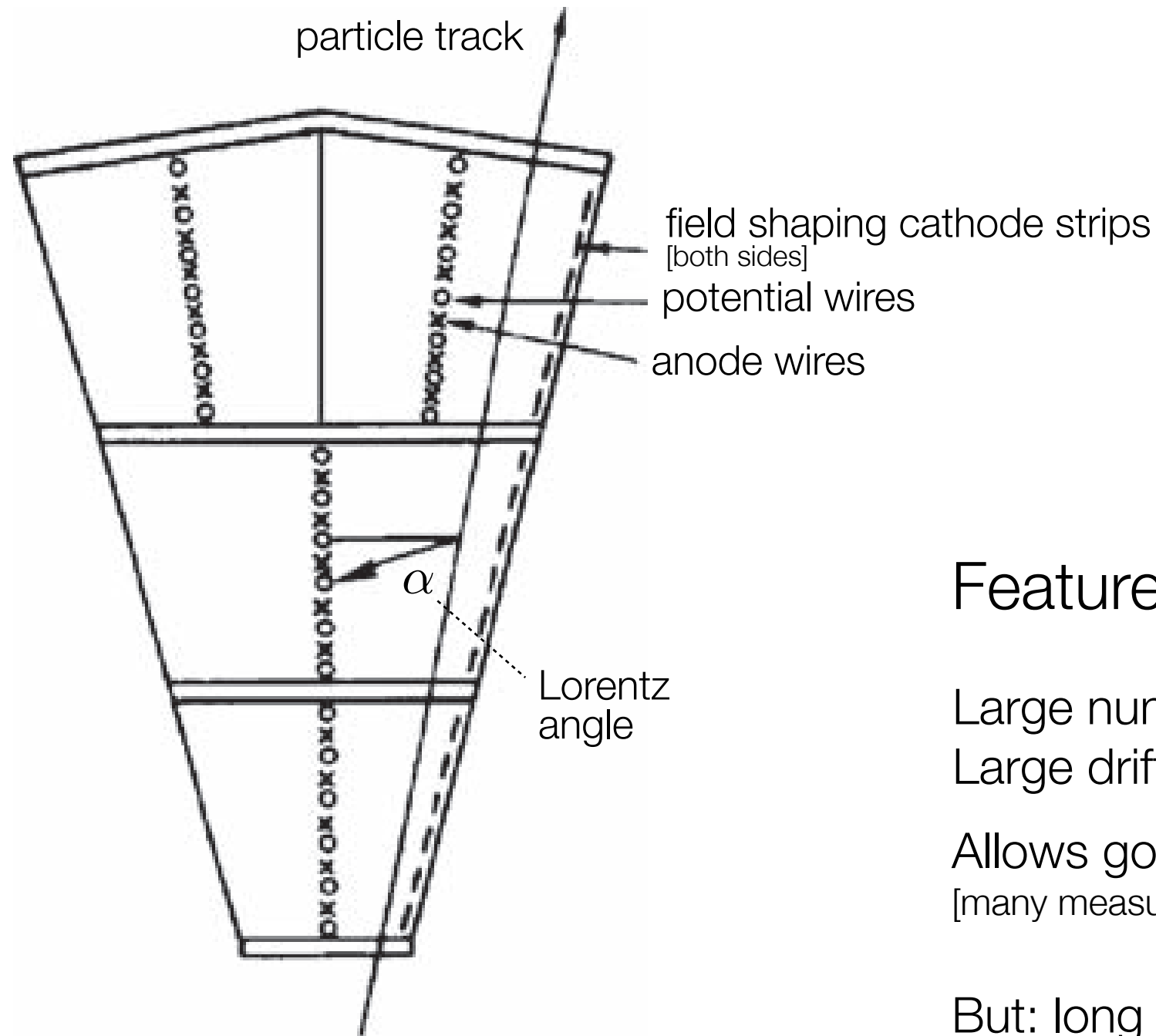
Avoid unnecessary contamination ...

Harmful are ...

- Halogens or halogen compounds
- Silicon compounds
- Carbonates, halocarbons
- Polymers
- Oil, fat ...
-



Jet drift chambers



Features:

Large number of sense wires

Large drift cells

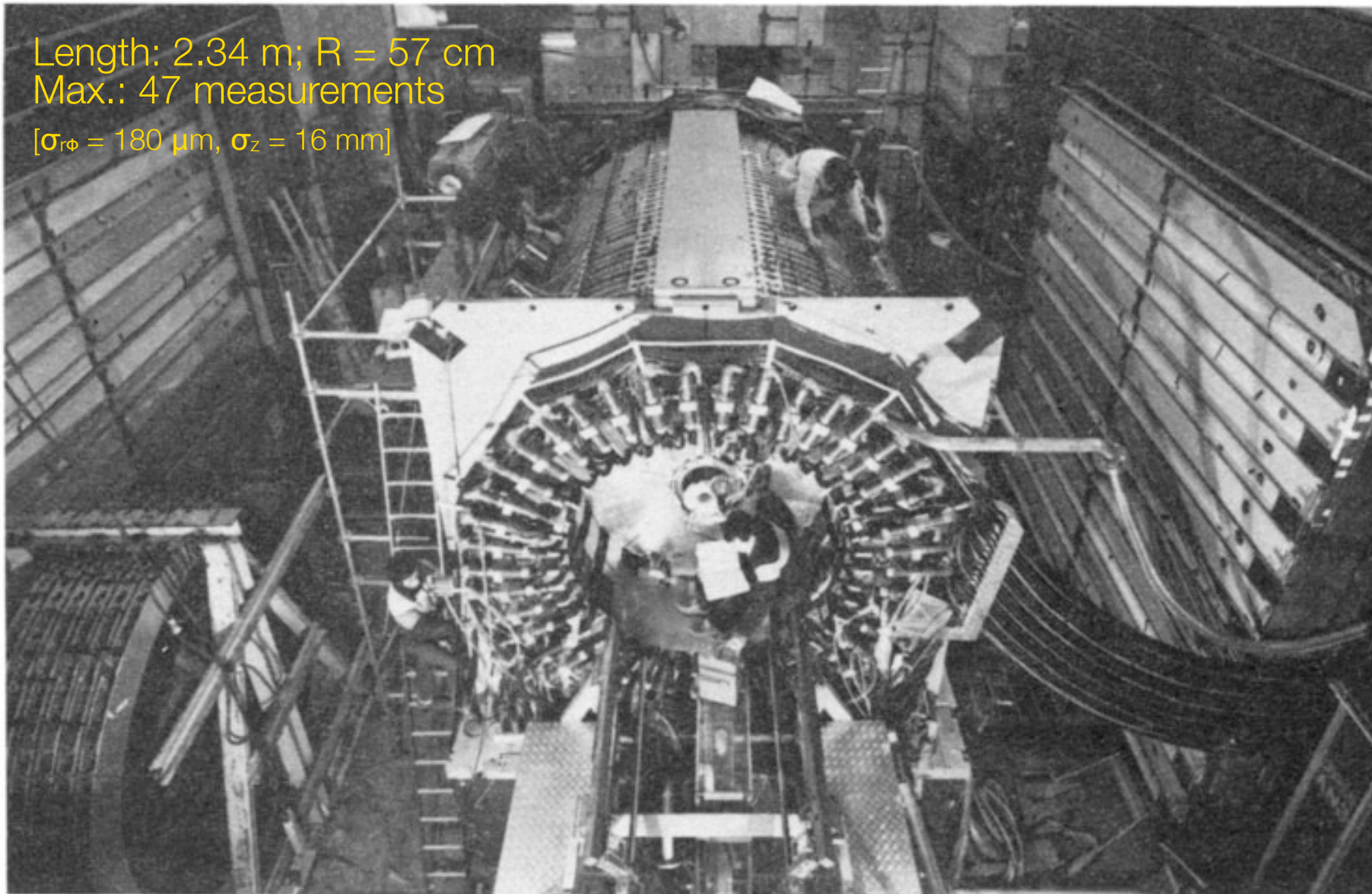
Allows good dE/dx determination
[many measurements]

But: long drift times ...

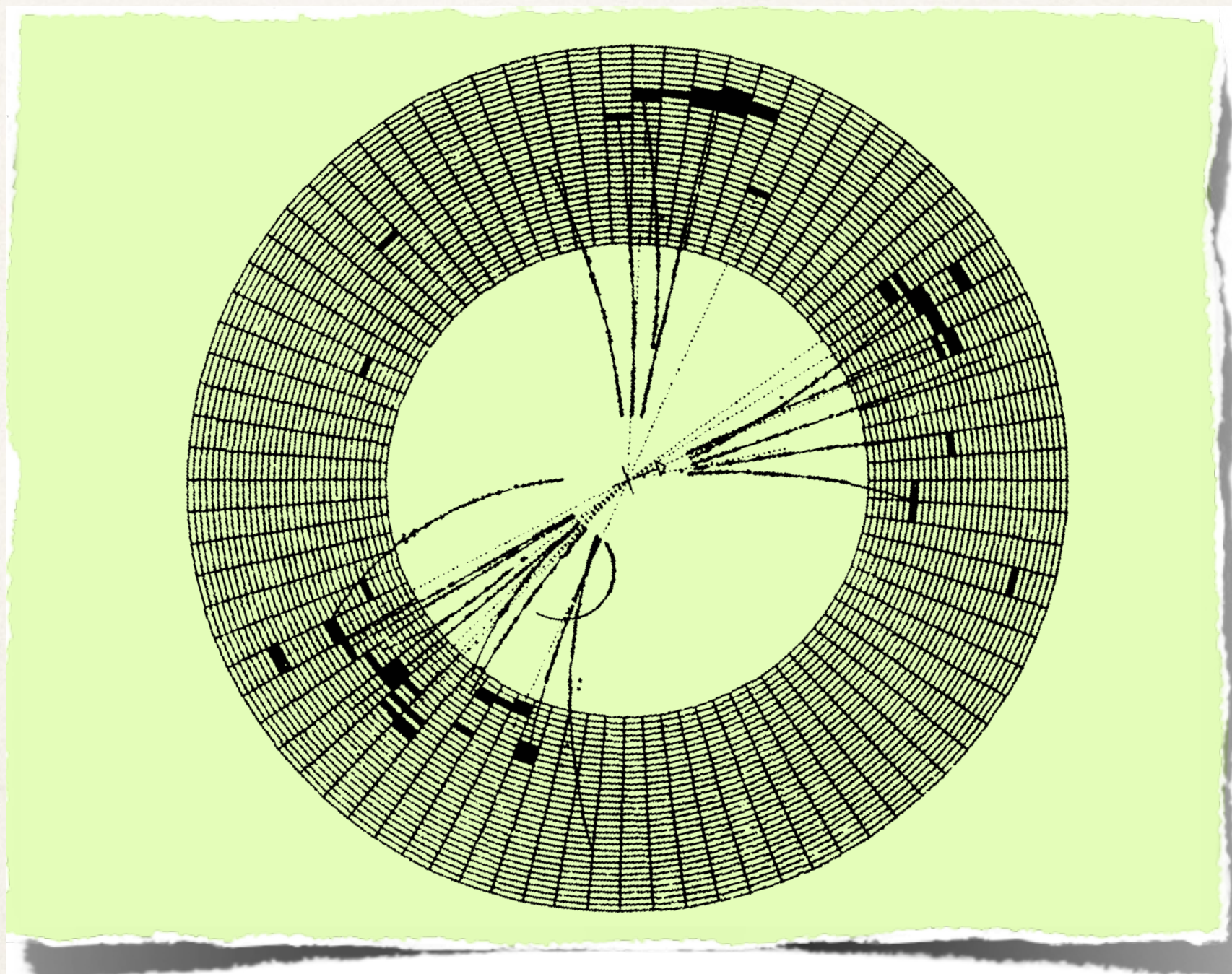
Jet drift chambers

Jet Drift Chambers – JADE

Length: 2.34 m; R = 57 cm
Max.: 47 measurements
[$\sigma_{r\phi} = 180 \mu\text{m}$, $\sigma_z = 16 \text{ mm}$]



Jet drift chambers



<http://cern.ch/amoraes>

Jet drift chambers

