



Estados de base: $\{|u_n\rangle\}$, $\langle u_m | u_n \rangle = \delta_{m,n}$

$$|\Psi(t)\rangle = \sum_n c_n(t) |u_n\rangle, \quad \sum_n |c_n(t)|^2 = 1$$

Projeção:

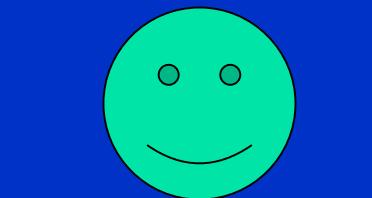
$$\langle u_p | \Psi(t) \rangle = c_p \langle u_p | u_p \rangle = c_p \quad \text{ou} \quad \langle \Psi(t) | u_p \rangle = c_p^* \langle u_p | u_p \rangle = c_p^*$$

Representação Matricial:

$$\langle u_n | A | u_p \rangle = A_{np} \quad \langle A \rangle(t) = \langle \Psi(t) | A | \Psi(t) \rangle = \underbrace{\sum_{np} c_n^* c_p A_{np}}$$

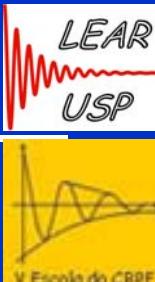
Equação de Schrödinger:

$$i\hbar \frac{d}{dt} |\Psi(t)\rangle = H(t) |\Psi(t)\rangle$$





Operador e Matriz Densidade



$$c_n^*(t) c_p(t) = [\langle u_p | \Psi(t) \rangle] [\langle \Psi(t) | u_n \rangle] = \underbrace{\langle u_p | \Psi(t) \rangle \langle \Psi(t) | u_n \rangle}_{= \langle u_p | \rho(t) | u_n \rangle} = \rho_{pn}(t)$$

$$\sum_n |c_n(t)|^2 = \sum_n \rho_{nn}(t) = Tr \rho(t) = 1$$

$$\langle A \rangle(t) = Tr \{ \rho(t) A \}$$

Equação de Liouville - von Neumann:

$$\frac{d}{dt} \rho(t) = \frac{1}{i\hbar} [\mathcal{H}(t), \rho(t)]$$



Evolução temporal do operador densidade



Se H é independente do tempo:

$$\frac{d}{dt} \rho(t) = \frac{1}{i\hbar} [H(t), \rho(t)]$$

Equação de Liouville - von Neumann

$$\rho(t) = e^{-iHt} \rho(0) e^{+iHt}, \quad U(t) = e^{-iHt}$$

$$\rho(t) = e^{-iH_n t_n} \dots \left[e^{-iH_3 t_3} \left[e^{-iH_2 t_2} \left(\underbrace{e^{-iH_1 t_1} \rho(0) e^{+iH_1 t_1}}_{\text{evento 1}} \right) e^{+iH_2 t_2} \right] e^{+iH_3 t_3} \dots \underbrace{e^{+iH_n t_n}}_{\text{n-ésimo evento}} \right]$$

evento 2
evento 3



Estado Puro vs. Estado Misto



Estado Puro:

$$|\Psi\rangle = \sum_n c_n |u_n\rangle \Rightarrow E_n$$



RMN: Mistura Estatística de Estados

$$|\Psi_1\rangle, |\Psi_2\rangle, \dots, |\Psi_k\rangle \Rightarrow p_1, p_2, \dots, p_k \Rightarrow E_1, E_2, \dots, E_k$$

$$p_1 + p_2 + \dots + p_k = 1$$



Populações e Coerências - Estado Misto

Qual é o significado físico dos elementos de matriz ρ_{np} em uma base: $\{|u_n\rangle\}$?

$$\rho_{nn} = \sum_k p_k [\rho_k]_{nn} = \sum_k p_k |c_n^{(k)}|^2 \Rightarrow \text{População do estado } |u_n\rangle$$

$$\rho_{np} = \sum_k p_k c_n^{(k)} c_p^{(k)*}$$

$c_n^{(k)} c_p^{(k)*} \Rightarrow$ Efeitos de Interferência entre os estados $|u_n\rangle$ e $|u_p\rangle$

$\rho_{np} \neq 0 \Rightarrow$ Coerência entre os estados $|u_n\rangle$ e $|u_p\rangle$



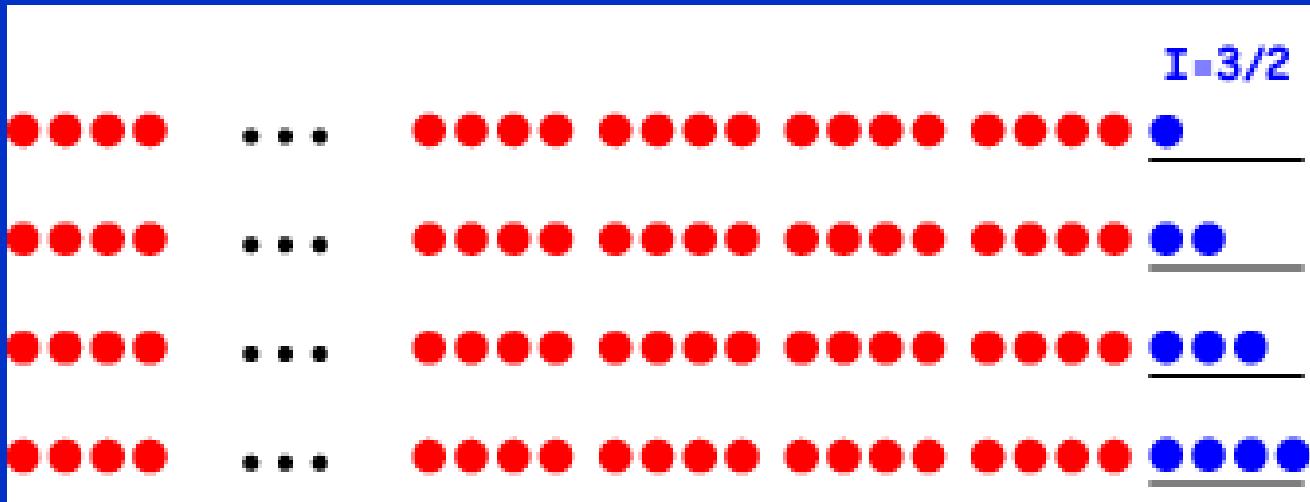
Sistema em equilíbrio térmico

Estatística: $\rho = \frac{e^{-H/kT}}{Z}, \quad Z = Tr\{e^{-H/kT}\}, \quad Tr \rho = 1$

$$\rho_{nn} = \frac{1}{Z} \langle u_n | e^{-H/kT} | u_n \rangle = \langle u_n | u_n \rangle \cdot \frac{e^{-E_n/kT}}{Z} = \frac{e^{-E_n/kT}}{Z} \quad (\text{Populações...})$$

$$\rho_{np} = \frac{1}{Z} \langle u_n | e^{-H/kT} | u_p \rangle = \langle u_n | u_p \rangle = 0 \quad (\text{Coerências...})$$

$$\rho_{nn} \approx \left(\hat{1} + \frac{n\hbar\omega_L}{kT} I_z + \dots \right)$$





Sistema em equilíbrio térmico: Zeeman

$$H = -\gamma \hbar I_z B_0 = \hbar \omega_L I_z$$

$$|H/kT| = |\hbar \omega_L I_z / kT| \ll 1$$

$$\rho = \frac{e^{-H/kT}}{Z} \approx \left(\hat{1} + \frac{\hbar \omega_L I_z}{kT} + \dots \right) \quad \frac{d}{dt} \rho(t) = \frac{1}{i\hbar} [H(t), \rho(t)]$$

Operador densidade parcial:

$$\Delta \rho \approx \frac{\hbar \omega_L}{kT} I_z = \alpha I_z \Rightarrow \boxed{\Delta \rho(0) = I_z} *$$

Spin 3/2: $\Delta \rho(0) = I_z = \begin{pmatrix} 3/2 & 0 & 0 & 0 \\ 0 & 1/2 & 0 & 0 \\ 0 & 0 & -1/2 & 0 \\ 0 & 0 & 0 & -3/2 \end{pmatrix}$



Sistema Laboratório:

$$\rho(t) = e^{-i\omega_L t I_z} I_\alpha e^{i\omega_L t I_z} \quad , \quad I_\alpha = I_x, I_y \text{ ou } I_z$$

Sistema Girante:

$$\rho_R(t) = e^{i\omega_R t I_z} \rho(t) e^{-i\omega_R t I_z} =$$

$$= e^{i\omega_R t I_z} \left(e^{-i\omega_L t I_z} I_\alpha e^{i\omega_L t I_z} \right) e^{-i\omega_R t I_z} =$$

$$= e^{i(\omega_R - \omega_L)t I_z} I_\alpha e^{-i(\omega_R - \omega_L)t I_z} = I_\alpha$$



Efeito do campo de radiofreqüência



$$B_1(t) = B_1 \cos(\omega_{RF} t + \phi)$$



$$H_{RF} = -\gamma \hbar I_\alpha B_{1\alpha}$$

$$e^{-i(\gamma B_{1\alpha})I_\alpha t} I_\alpha e^{i(\gamma B_{1\alpha})I_\alpha t} = e^{i\gamma B_{1\alpha} t} I_\alpha = e^{i\beta_\alpha} I_\alpha$$

$$\rho(t) = e^{-i\hbar(\gamma B_{1\alpha})I_x t} I_z e^{i\hbar(\gamma B_{1\alpha})I_\alpha t} = e^{i\gamma B_{1\alpha} t} I_z$$

$$\rho_{\pi/2,x}(t) = e^{i(\pi/2)_x} I_z = I_z \cos(\pi/2) - I_y \sin(\pi/2) = I_y$$

$$\rho_{\pi,y}(t) = e^{i(\pi)_y} I_z = I_z \cos(\pi) + I_x \sin(\pi) = -I_z$$



Sinal de RMN - Zeeman

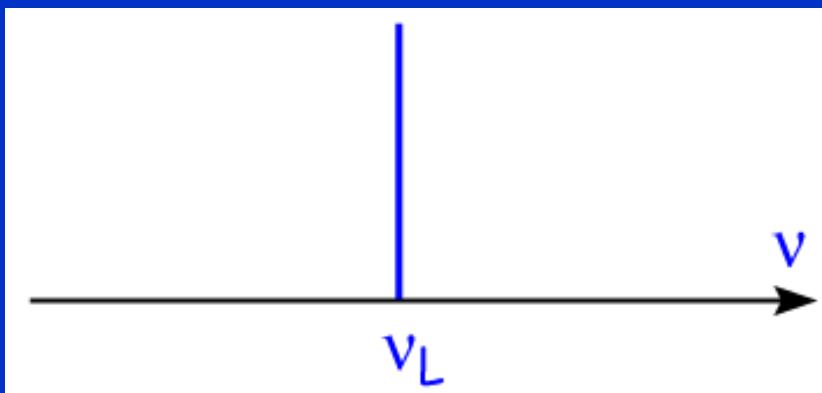


$$\rho(t) = e^{-i(\gamma B_0) I_z t} I_x e^{i(\gamma B_0) I_z t} = e^{i\gamma B_0 t} I_x =$$

$$= I_x \cos(\omega_L t) + I_y \sin(\omega_L t)$$

$$\vec{\mu} = \gamma \hbar \vec{I} \Rightarrow \vec{M} = \sum_i \langle \vec{\mu}_i \rangle$$

$$\langle M_{xy}(t) \rangle = \langle M_{xy,0} \rangle \exp(i\omega_L t)$$





Sinal de RMN - Quadrupolar - pulsos não seletivos

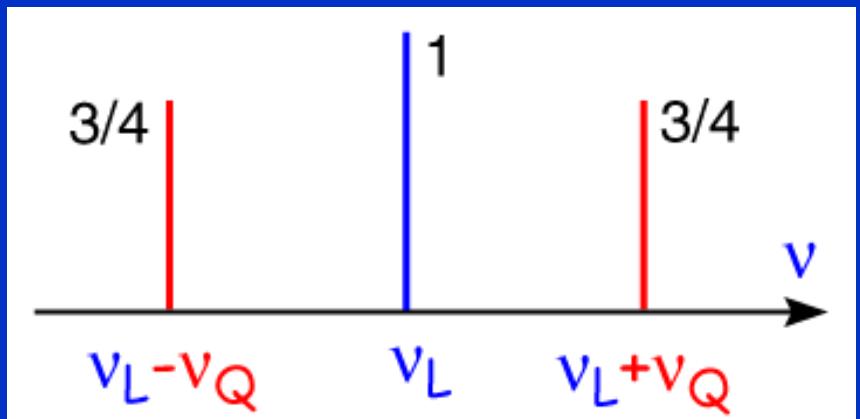
$$H_Q = \frac{1}{6} \omega_Q \left[3I_z^2 - I(I+1)\hat{1} \right] = \omega_Q \begin{bmatrix} 1/2 & 0 & 0 & 0 \\ 0 & -1/2 & 0 & 0 \\ 0 & 0 & -1/2 & 0 \\ 0 & 0 & 0 & 1/2 \end{bmatrix}$$



$$\rho(t) = e^{-i\omega_Q t} I_x e^{i\omega_Q t} = \begin{bmatrix} 0 & \sqrt{3/4}e^{-i\omega_Q t} & 0 & 0 \\ \sqrt{3/4}e^{i\omega_Q t} & 0 & 1 & 0 \\ 0 & 1 & 0 & \sqrt{3/4}e^{i\omega_Q t} \\ 0 & 0 & \sqrt{3/4}e^{-i\omega_Q t} & 0 \end{bmatrix}$$

$$\langle I_x(t) + iI_y(t) \rangle = Tr[\rho(t)(I_x + iI_y)] = 3\cos(\omega_Q t) + 2\cos(0.t)$$

↑
FID



TF



Sinal de RMN - Quadrupolar - pulsos seletivos

$$P_{\alpha}^{01}(\theta) = \begin{pmatrix} \cos\left(\frac{\sqrt{3}\theta}{2}\right)e^{2i\omega_q t_p} & i\sin\left(\frac{\sqrt{3}\theta}{2}\right)e^{-i\left(2\omega_q t_p + \frac{\pi}{2}\alpha\right)} & 0 & 0 \\ i\sin\left(\frac{\sqrt{3}\theta}{2}\right)e^{-i\left(2\omega_q t_p - \frac{\pi}{2}\alpha\right)} & \cos\left(\frac{\sqrt{3}\theta}{2}\right)e^{2i\omega_q t_p} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & e^{-i4\omega_q t_p} \end{pmatrix}$$



$$P_{\alpha}^{12}(\theta) = \begin{pmatrix} e^{-i\omega_q t_p} & 0 & 0 & 0 \\ 0 & \cos\left(\frac{\theta}{2}\right)e^{i\omega_q t_p} & i\sin\left(\frac{\theta}{2}\right)e^{-i\left(\omega_q t_p + \frac{\pi}{2}\alpha\right)} & 0 \\ 0 & i\sin\left(\frac{\theta}{2}\right)e^{-i\left(\omega_q t_p - \frac{\pi}{2}\alpha\right)} & \cos\left(\frac{\theta}{2}\right)e^{i\omega_q t_p} & 0 \\ 0 & 0 & 0 & e^{-i\omega_q t_p} \end{pmatrix}$$

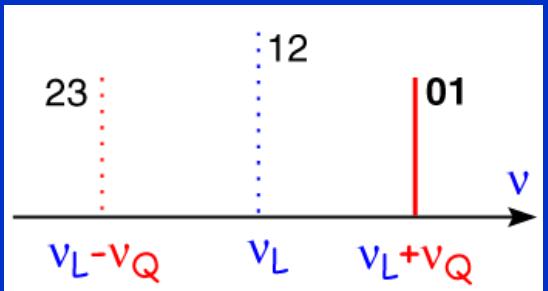
$$P_{\alpha}^{23}(\theta) = \begin{pmatrix} e^{-2i\omega_q t_p} & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & \cos\left(\frac{\sqrt{3}\theta}{2}\right)e^{2i\omega_q t_p} & i\sin\left(\frac{\sqrt{3}\theta}{2}\right)e^{-i\left(2\omega_q t_p + \frac{\pi}{2}\alpha\right)} \\ 0 & 0 & i\sin\left(\frac{\sqrt{3}\theta}{2}\right)e^{-i\left(2\omega_q t_p - \frac{\pi}{2}\alpha\right)} & \cos\left(\frac{\sqrt{3}\theta}{2}\right)e^{2i\omega_q t_p} \end{pmatrix}$$



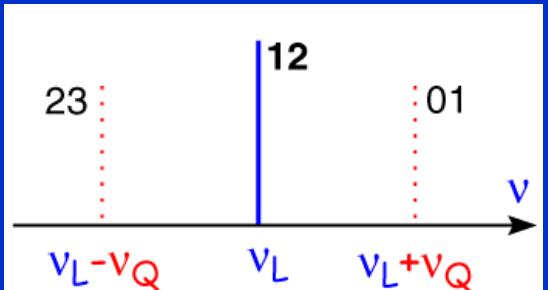
Sinal de RMN - Quadrupolar - pulsos seletivos



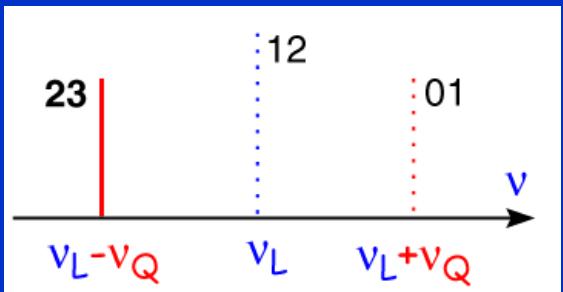
$$P_y^{01} \Rightarrow \langle I_x(t) + iI_y(t) \rangle = Tr[\rho(t)(I_x + iI_y)] = \sqrt{3/4}e^{+i\omega_Q t}$$



$$P_y^{12} \Rightarrow \langle I_x(t) + iI_y(t) \rangle = Tr[\rho(t)(I_x + iI_y)] = 1$$



$$P_y^{23} \Rightarrow \langle I_x(t) + iI_y(t) \rangle = Tr[\rho(t)(I_x + iI_y)] = \sqrt{3/4}e^{-i\omega_Q t}$$





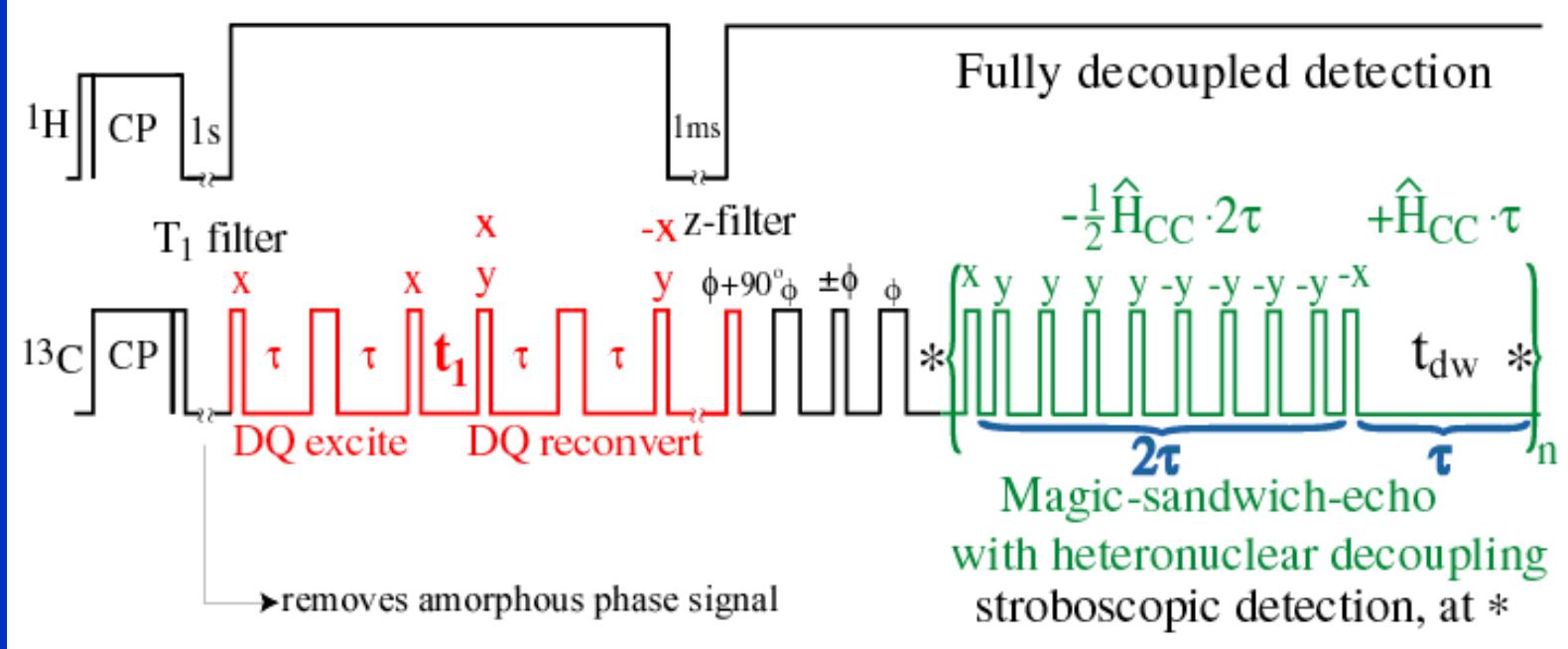
- Mecânica Quântica
- Espectrômetro de RMN
- Química
- Física do problema...

- Pulso de RF
- Tripla ressonância
- Campos magnéticos
- Rotação da amostra
- Manipulação química
- Temperatura

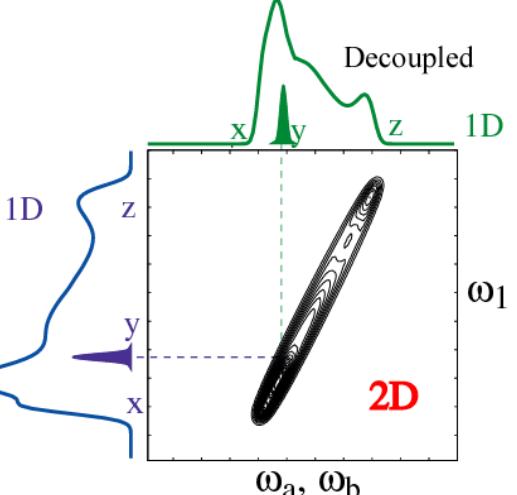
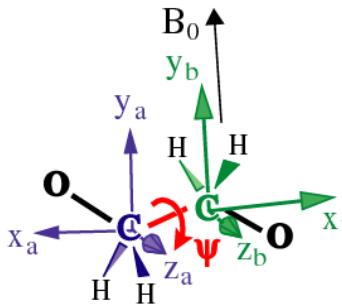


Estudos de conformação por RMN molecular...

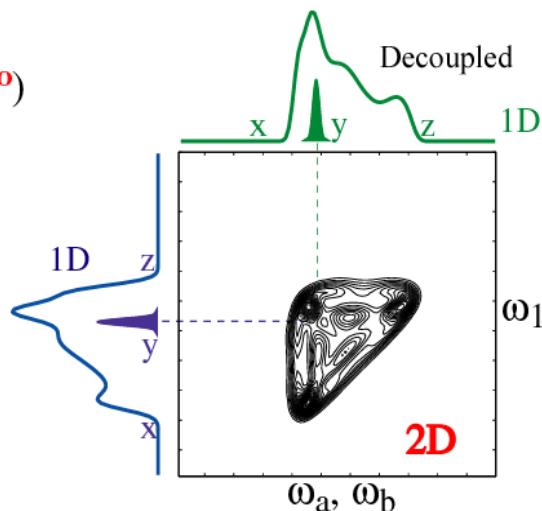
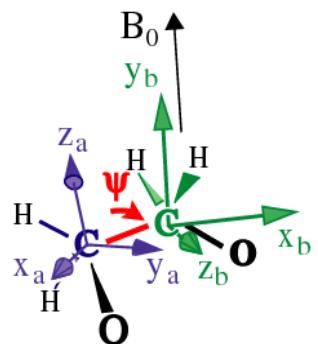
Double- -Quantum NMR



trans ($\psi = 170^\circ$)

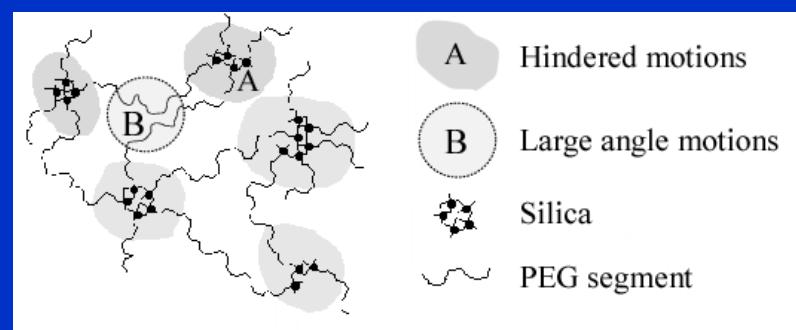
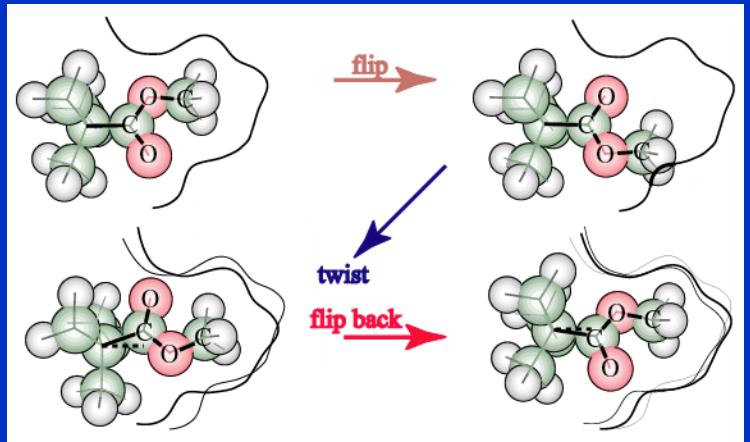


gauche ($\psi = \pm 60^\circ$)

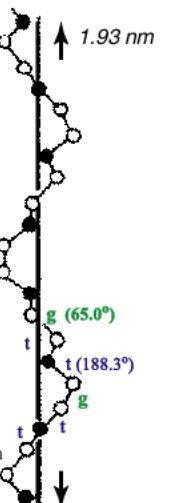




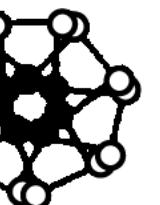
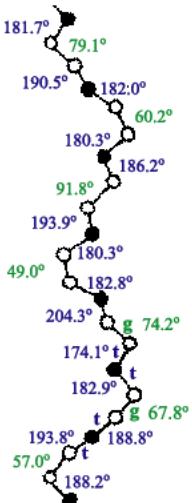
Exemplos... RMN...



Conformation of Crystalline PEO
Semicrystalline poly(ethylene oxide), PEO, $[CH_2-CH_2-O-]_n$
Chain structures suggested based on X-ray fiber diffraction:
Ideal 7_2 helix Distorted 7_2 helix



side view



top view



Takahashi & Tadokoro, 1973
 $O-C-C-O: 68^\circ \pm 14^\circ$

lokoro et al., 1964
 $C-C-O: \psi = 65^\circ$



Matrizes exponenciais - Diagonalização

Sendo D diagonal: $e^D = \sum_{N=0}^{\infty} \frac{D^N}{N!} = \begin{bmatrix} e^{D_{11}} & 0 & \dots & 0 \\ 0 & e^{D_{22}} & \dots & 0 \\ \dots & \dots & \dots & \dots \\ 0 & 0 & \dots & e^{D_{ii}} \end{bmatrix}$

Sendo A não Diagonal \rightarrow diagonalizar A :

$$\left(R_D^{-1} A R_D \right)_{mn} = \delta_{mn} a_n \Rightarrow e^{R_D^{-1} A R_D} = \begin{bmatrix} e^{a_{11}} & 0 & \dots & 0 \\ 0 & e^{a_{22}} & \dots & 0 \\ \dots & \dots & \dots & \dots \\ 0 & 0 & \dots & e^{a_{ii}} \end{bmatrix}$$

$$e^{R_D^{-1} A R_D} = R_D^{-1} e^A R_D \Rightarrow$$

$$e^A = R_D \left[R_D^{-1} e^A R_D \right] R_D^{-1} = R_D \begin{bmatrix} e^{a_{11}} & 0 & \dots & 0 \\ 0 & e^{a_{22}} & \dots & 0 \\ \dots & \dots & \dots & \dots \\ 0 & 0 & \dots & e^{a_{ii}} \end{bmatrix} R_D^{-1}$$