

Black Hole Remnants

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-Mimetic gravity

$$S = \frac{1}{2} \int d^4x \sqrt{-g} \left(-\frac{1}{8\pi G} R + \lambda (g^{\mu\nu} \phi_{,\mu} \phi_{,\nu} - 1) \right),$$

- *Constraint*

$$g^{\mu\nu} \phi_{,\mu} \phi_{,\nu} = 1,$$

- *Synchronous coordinate system*

$$ds^2 = dt^2 - \gamma_{ik} dx^i dx^k$$

$$\phi = t$$

Extrinsic curvature of hypersurfaces $t = \text{const}$

$$\kappa_{ik} = \frac{1}{2} \frac{\partial}{\partial t} \gamma_{ik} = -\phi_{;ik}$$

and

$$\square\phi = g^{\alpha\beta} \phi_{;\alpha\beta} = \gamma^{ik} \kappa_{ik} = \kappa = \frac{\partial}{\partial t} \ln \sqrt{\gamma},$$

• **Action for modified mimetic gravity**

$$S = \frac{1}{2} \int d^4x \sqrt{-g} (-f(\square\phi) R - 2\Lambda(\square\phi) + \lambda(g^{\mu\nu} \phi_{,\mu} \phi_{,\nu} - 1) + 2\mathcal{L}_m)$$

where

$$f(\square\phi) = \frac{1}{8\pi G(\square\phi)}$$

"Asymptotic freedom"

$$G \propto \frac{1}{f} \rightarrow 0 \quad \text{when} \quad \square\phi = \kappa \rightarrow \kappa_0$$

$$f = \frac{1}{1 - (\kappa^2/\kappa_0^2)},$$

- **Friedmann Universe**

$$ds^2 = dt^2 - a^2(t) \delta_{ik} dx^i dx^k$$

$$\frac{1}{3} \kappa^2 \left(\frac{1 + 2(\kappa^2/\kappa_0^2)}{1 - (\kappa^2/\kappa_0^2)} \right) = \varepsilon \quad \text{where} \quad \kappa = 3 \frac{\dot{a}}{a}$$

$$a \propto t^{\frac{2}{3(1+w)}} \quad \text{for} \quad \kappa^2 / \kappa_0^2 \ll 1$$

$$a \propto \exp\left(-\frac{\kappa_0 t}{3}\right) \quad \text{when} \quad \kappa^2 / \kappa_0^2 \rightarrow \mathbf{1}$$

● **Kasner Universe**

$$\gamma_{ik} = \gamma_{(i)}(t) \delta_{ik} \quad \text{where} \quad \gamma_{(i)} = \left(\frac{3}{2} \bar{\lambda}^2\right)^{1/3} t^{2p_i},$$

$$\text{and } p_1 + p_2 + p_3 = 1, \quad p_1^2 + p_2^2 + p_3^2 = 1 \quad \text{for } \kappa^2 / \kappa_0^2 \ll 1$$

$$\gamma_{(1)} = \gamma_{(2)} = \gamma_{(3)} = \gamma^{1/3} \propto \exp\left(-\frac{2}{3} \kappa_0 t\right) \quad \text{when } \kappa^2 / \kappa_0^2 \rightarrow \mathbf{1}$$

Quantum fluctuations

$$\delta h(l) \simeq \frac{1}{\sqrt{fl}} \simeq \frac{\sqrt{G(\kappa)}}{l}$$

Hence, quantum fluctuations in a given physical scale $l \ll \kappa_0^{-1}$ vanish as $\kappa \rightarrow \kappa_0$ and correspondingly $G(\kappa) \rightarrow 0$.

Black Hole Remnants

$$ds^2 = (1 - a^2(r)) dt^2 - \frac{dr^2}{(1 - a^2(r))} - r^2 d\Omega^2$$

where $a^2(r) = r_g/r$ *for* **BH**
and $a^2 = (Hr)^2$ *for* **de Sitter**.

The Lemaître coordinates

$$T = t + \int \frac{a}{1 - a^2} dr, \quad R = t + \int \frac{dr}{a(1 - a^2)}$$

$$ds^2 = dT^2 - a^2(x) dR^2 - b^2(x) d\Omega^2 \quad \textit{where} \quad x \equiv R - T$$

BH

$$ds^2 = dT^2 - (x/x_+)^{-2/3} dR^2 - (x/x_+)^{4/3} r_g^2 d\Omega^2,$$

and it is regular at the horizon $x = x_+ = 4M/3$.
region $x > 0$ covers both interior and exterior of the black

de Sitter

$$ds^2 = dT^2 - \exp(2H(x - x_-)) (dR^2 + H^{-2}d\Omega^2)$$

where x_- is a constant of integration in (4) and the de Sitter horizon occurs at $x = x_-$. The region $x < x_-$ corresponds to the patch of size $r = H^{-1}$ covered by static coordinates, which on larger scales do not exist.

Asymptotically free mimetic gravity with

$$f(\tilde{\kappa}) = \frac{1 + 3\tilde{\kappa}^2}{(1 + \tilde{\kappa}^2)(1 - \tilde{\kappa}^2)^2}, \quad \text{where } \tilde{\kappa} \equiv \kappa/\kappa_0$$

Exact solution

$$ds^2 = dT^2 - a^2(x) dR^2 - b^2(x) d\Omega^2$$

where

$$a^3(\tilde{\kappa}) = \frac{4M\kappa_0}{3} |\tilde{\kappa}| (1 - \tilde{\kappa}^4) \left(\frac{1 + \tilde{\kappa}^2}{1 + 3\tilde{\kappa}^2} \right)^2,$$

$$b^3(\tilde{\kappa}) = \frac{9M}{2\kappa_0^2 \tilde{\kappa}^2} (1 - \tilde{\kappa}^2) (1 + 3\tilde{\kappa}^2).$$

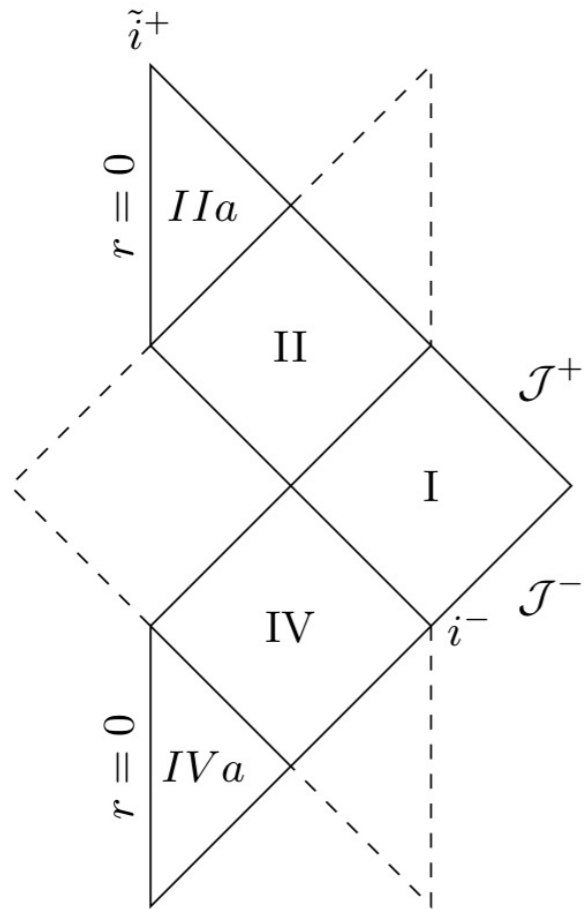
$$-\kappa_0 x = \frac{1}{\tilde{\kappa}} + 2 (\arctan \tilde{\kappa} - \tanh^{-1} \tilde{\kappa}).$$

For $x \rightarrow +\infty$ *we have BH*

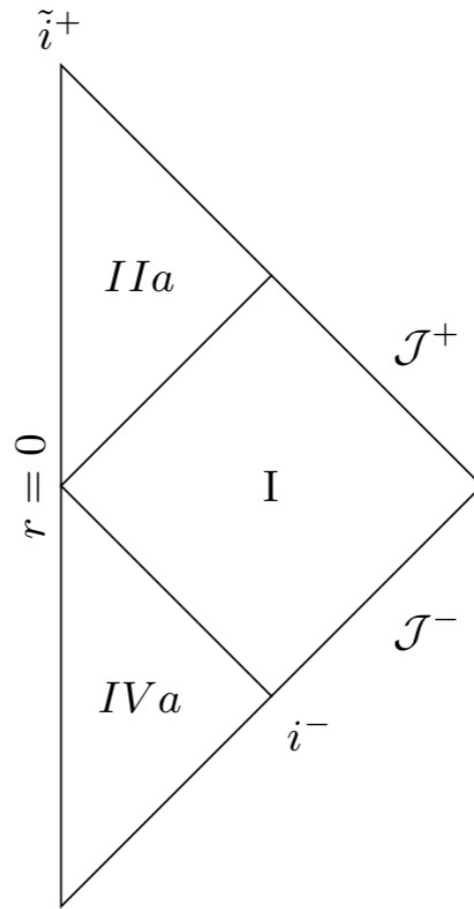
When $x \rightarrow -\infty$ *we obtain static de Sitter patch*

Horizons exist only if

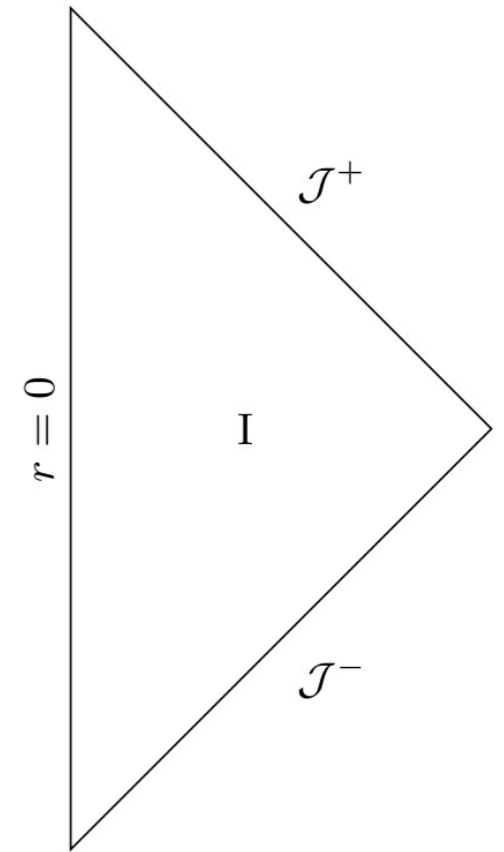
$$M \geq M_{\min} = \frac{5^{5/2}}{18\kappa_0}.$$



$M > M_0$



$M = M_0$



$M < M_0$

Near horizon metric for $M = M_{\min}$

$$ds^2 \approx \frac{10}{7} \left(1 - \frac{r}{r_*}\right)^2 dt^2 - \frac{dr^2}{\frac{10}{7} \left(1 - \frac{r}{r_*}\right)^2} - r^2 d\Omega^2$$

where $r_* = \left(\frac{6}{5}\right)^2 M_{\min}$

Black hole thermodynamics

$$T_H = \frac{g_s}{2\pi} = \frac{\kappa_0}{6\pi} |\tilde{\kappa}_+| \frac{1 - 5\tilde{\kappa}_+^2}{1 + 3\tilde{\kappa}_+^2}, \quad M = \frac{3}{4\kappa_0 |\tilde{\kappa}_+| (1 - \tilde{\kappa}_+^4)} \left(\frac{1 + 3\tilde{\kappa}_+^2}{1 + \tilde{\kappa}_+^2} \right)^2.$$

