Asymptotically Safe Quantum Gravity

Antonio D. Pereira

Universidade Federal Fluminense & Radboud University Nijmegen

XVIII Brazilian School of Cosmology and Gravitation









Universidade Federal Fluminense





Recorded Lectures...

Recorded Lectures...

...in Portuguese

Recorded Lectures...

... in Portuguese



Cursos - Introdução ao grupo de renormalização funcional (Antônio D. Pereira)

20 vídeos • 919 visualizações • Última atualização em 4 de jun. de 2020

=+ 次 谷 …

Este curso foi dado em 2016 para alunos de pós-

MOSTRAR MAIS





Cursos - Gravidade quântica assintoticamente segura

23 vídeos • 571 visualizações • Última atualização em 6 de jun. de 2020

Dr. Anderson Tomaz e Dr. Gustavo P. de Brito

Curso oferecido pelo Dr. Anderson Tomaz e o estudante de doutorado Gustavo P. de Brito nas Atividades Formativas de Verão no ano de 2018 (no CBPF). O curso consiste em uma introdução à abordagem para a quantização da interação gravitacional conhecida como "Gravidade quântica assintoticamente segura" (Asymptotically safe quantum gravity).

Starting point: single scalar field

Starting point: single scalar field

$$Z[J] = \int_{\Lambda} \mathscr{D}\phi \,\mathrm{e}^{-S[\phi] + \int \mathrm{d}^d x \, J(x)\phi(x)}$$

generating functional of correlation functions

Starting point: single scalar field

$$Z[J] = \int_{\Lambda} \mathscr{D}\phi \,\mathrm{e}^{-S[\phi] + \int \mathrm{d}^d x \, J(x)\phi(x)}$$

generating functional of correlation functions

 $W[J] = \ln Z[J]$

generating functional of connected correlation functions

Starting point: single scalar field

$$Z[J] = \int_{\Lambda} \mathscr{D}\phi \,\mathrm{e}^{-S[\phi] + \int \mathrm{d}^d x \, J(x)\phi(x)}$$

generating functional of correlation functions

 $W[J] = \ln Z[J]$ generating functional of *connected* correlation functions

$$\langle \phi(x_1)...\phi(x_n) \rangle = \frac{\delta^n W[J]}{\delta J(x_1)...\delta J(x_n)} \bigg|_{J=0}$$

Starting point: single scalar field

$$Z[J] = \int_{\Lambda} \mathscr{D}\phi \,\mathrm{e}^{-S[\phi] + \int \mathrm{d}^d x \, J(x)\phi(x)}$$

generating functional of correlation functions

 $W[J] = \ln Z[J]$

generating functional of *connected* correlation functions

$$\langle \phi(x_1)...\phi(x_n) \rangle = \frac{\delta^n W[J]}{\delta J(x_1)...\delta J(x_n)} \bigg|_{J=0}$$

$$\varphi(x) \equiv \langle \phi(x) \rangle_J = \frac{\delta W[J]}{\delta J(x)}$$

Starting point: single scalar field

$$Z[J] = \int_{\Lambda} \mathscr{D}\phi \,\mathrm{e}^{-S[\phi] + \int \mathrm{d}^d x \, J(x)\phi(x)}$$

generating functional of correlation functions

 $W[J] = \ln Z[J]$

generating functional of *connected* correlation functions

$$\langle \phi(x_1)...\phi(x_n) \rangle = \frac{\delta^n W[J]}{\delta J(x_1)...\delta J(x_n)} \bigg|_{J=0}$$

$$\varphi(x) \equiv \langle \phi(x) \rangle_J = \frac{\delta W[J]}{\delta J(x)}$$

$$\Gamma[\varphi] = - W[J_{\varphi}] + \int d^d x \ J_{\varphi} \ \varphi$$

Starting point: single scalar field

$$Z[J] = \int_{\Lambda} \mathscr{D}\phi \,\mathrm{e}^{-S[\phi] + \int \mathrm{d}^d x \, J(x)\phi(x)}$$

generating functional of correlation functions

 $W[J] = \ln Z[J]$

generating functional of *connected* correlation functions

$$\langle \phi(x_1)...\phi(x_n) \rangle = \frac{\delta^n W[J]}{\delta J(x_1)...\delta J(x_n)} \bigg|_{J=0}$$

$$\varphi(x) \equiv \langle \phi(x) \rangle_J = \frac{\delta W[J]}{\delta J(x)}$$

$$\Gamma[\varphi] = -W[J_{\varphi}] + \int \mathrm{d}^d x \ J_{\varphi} \ \varphi$$

generating functional of one-particle irreducible correlation functions

$$Z[J] = \int_{\Lambda} \mathscr{D}\phi \,\mathrm{e}^{-S[\phi] + \int \mathrm{d}^d x \, J(x)\phi(x)}$$

regulator "action":

$$\Delta S_k = \frac{1}{2} \int \mathrm{d}^d x \, \phi(x) \, \mathcal{R}_k(-\partial^2) \phi(x)$$

regulator "action":

$$\Delta S_k = \frac{1}{2} \int d^d x \, \phi(x) \, \mathcal{R}_k(-\partial^2) \phi(x)$$
gives a (large) mass to
field modes with

momentum lower than k









finally

$$\Gamma_k[\varphi] \equiv \overline{\Gamma}_k[\varphi] - \frac{1}{2} \int d^d x \ \varphi(x) \mathcal{R}_k(-\partial^2) \varphi(x)$$

effective average action (EAA)

interpolates between full effective action and the "classical" one



interpolates between full effective action and the "classical" one

satisfies an exact flow equation



interpolates between full effective action and the "classical" one

satisfies an exact flow equation



$$\partial_t \Gamma_k = \frac{1}{2} \operatorname{STr} \left[\left(\Gamma_k^{(2)} + \mathscr{R}_k \right)^{-1} \partial_t \mathscr{R}_k \right]$$

 $\partial_t \equiv k \partial_k$

(exact) flow equation

Wetterich equation

interpolates between full effective action and the "classical" one

satisfies an exact flow equation



$$\partial_t \Gamma_k = \frac{1}{2} \operatorname{STr} \left[\left(\Gamma_k^{(2)} + \mathcal{R}_k \right)^{-1} \partial_t \mathcal{R}_k \right] \qquad \qquad \partial_t \equiv k \partial_k$$

(exact) flow equation

Wetterich equation

conversion of functional integral into functional differential equation





space of all functionals of the field which are compatible with the symmetries of the theory

(Infinitely many)

space of all functionals of the field which are compatible with the symmetries of the theory

(Infinitely many)

space of all functionals of the field which are compatible with the symmetries of the theory

the effective average action is expanded as

$$\Gamma_k[\varphi] = \sum_i \bar{g}_i(k) \mathcal{O}_i[\varphi]$$

(Infinitely many)

space of all functionals of the field which are compatible with the symmetries of the theory

the effective average action is expanded as

$$\Gamma_k[\varphi] = \sum_i \bar{g}_i(k) \mathcal{O}_i[\varphi]$$

$$\partial_t \Gamma_k[\varphi] = \sum_i \left(\partial_t \bar{g}_i(k) \right) \mathcal{O}_i[\varphi]$$

$$\bar{g}_i = k^{d_i} g_i$$
$$\partial_t \bar{g}_i = k^{d_i} (d_i g_i + \beta_i)$$

$$\beta_i = -d_i g_i$$

 $+k^{-d_i}\partial_t \bar{g}_i$

(Infinitely many)

space of all functionals of the field which are compatible with the symmetries of the theory

the effective average action is expanded as

$$\Gamma_{k}[\varphi] = \sum_{i} \bar{g}_{i}(k) \mathcal{O}_{i}[\varphi]$$

$$\bar{g}_{i} = k^{d_{i}}g_{i}$$

$$\partial_{t}\Gamma_{k}[\varphi] = \sum_{i} (\partial_{t}\bar{g}_{i}(k)) \mathcal{O}_{i}[\varphi]$$

$$\bar{g}_{i} = k^{d_{i}}(d_{i}g_{i} + \beta_{i})$$

$$\beta_{i} = -d_{i}g_{i} + k^{-d_{i}}\partial_{t}\bar{g}_{i}$$

$$\beta_{i} = -d_{i}g_{i} + k^{-d_{i}}\partial_{t}\bar{g}_{i}$$
extraction of beta functions
suitable projection
rule for the Wetterich
equation
equation

(Infinitely many)

space of all functionals of the field which are compatible with the symmetries of the theory

the effective average action is expanded as

$$\Gamma_{k}[\varphi] = \sum_{i} \bar{g}_{i}(k) \mathcal{O}_{i}[\varphi]$$

$$\bar{g}_{i} = k^{d_{i}}g_{i}$$

$$\partial_{i}\bar{g}_{i} = k^{d_{i}}(d_{i}g_{i} + \beta_{i})$$

$$\beta_{i} = -d_{i}g_{i} + k^{-d_{i}}\partial_{i}\bar{g}_{i}$$

$$\beta_{i} = -d_{i}g_{i} + k^{-d_{i}}\partial_{i}\bar{g}_{i}$$
extraction of beta functions
suitable projection
rule for the Wetterich
equation
extraction of beta functions

Approximations are necessary, but we don't need to use a perturbative scheme!

Looking for fixed points:

Looking for fixed points:

$$\beta_i(\mathbf{g}^*) = 0, \ i = 1,...,\infty$$

 $\mathbf{g}^* = (g_1^*, \dots, g_\infty^*)$

Looking for fixed points:

$$\beta_i(\mathbf{g}^*) = 0, \ i = 1,...,\infty$$

 $\mathbf{g}^* = (g_1^*, \dots, g_\infty^*)$






diagonalize

 $\partial_t z_i = \lambda_i \, z_i$



diagonalize



Looking for fixed points:

$$\beta_i(\mathbf{g}^*) = 0, \ i = 1,...,\infty$$

 $\mathbf{g}^* = (g_1^*, \dots, g_\infty^*)$

$$\partial_t(g_i - g_i^*) = \sum_j \frac{\partial \beta_i}{\partial g_j}(g_j - g_j^*)$$

diagonalize













Asymptotic Safety:

Existence of a renormalization-group fixed point;

Fixed point features finitely many relevant directions;



Asymptotic Safety:

Existence of a renormalization-group fixed point;

Fixed point features finitely many relevant directions;



Approximations are necessary - Truncations of the effective average action

Approximations are necessary - Truncations of the effective average action

$$\Gamma_k[\varphi] = \sum_{i=1}^N \bar{g}_i(k) \mathcal{O}_i[\varphi]$$

Approximations are necessary - Truncations of the effective average action



Homework: Scalar Field Theory

$$\Gamma_{k}[\varphi] = \int d^{d}x \left(\frac{Z_{\varphi}}{2} \partial^{\mu} \varphi \partial_{\mu} \varphi + V_{k}(\varphi^{2}) \right)$$
$$V_{k}(\varphi^{2}) = \sum_{i=1}^{N} Z_{\varphi}^{i} \frac{\bar{g}_{2i}}{(2i)!} \varphi^{2i}$$



Asymptotically Safe Quantum Gravity

Antonio D. Pereira

Universidade Federal Fluminense & Radboud University Nijmegen

XVIII Brazilian School of Cosmology and Gravitation









Universidade Federal Fluminense





Asymptotically Safe Quantum Gravity

- the technical side -



Quantum Gravity and the Functional Renormalization Group

The Road towards Asymptotic Safety

MARTIN REUTER AND FRANK SAUERESSIG

CAMBRIDGE MONOGRAPHS ON MATHEMATICAL PHYSICS

Asymptotically Safe Quantum Gravity

- the technical side -

Asymptotically Safe Quantum Gravity

- the technical side -





No background to set a scale: background field method

$$g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}$$



No background to set a scale: background field method

$$g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}$$

The spectrum of the Laplacian of the background metric defines a scale

background independence is
encoded in split symmetry $\bar{g}_{\mu\nu} \rightarrow \bar{g}_{\mu\nu} + \epsilon_{\mu\nu}$
 $h_{\mu\nu} \rightarrow h_{\mu\nu} - \epsilon_{\mu\nu}$



No background to set a scale: background field method

$$g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}$$

The spectrum of the Laplacian of the background metric defines a scale

background independence is
encoded in split symmetry $\bar{g}_{\mu\nu} \rightarrow \bar{g}_{\mu\nu} + \epsilon_{\mu\nu}$
 $h_{\mu\nu} \rightarrow h_{\mu\nu} - \epsilon_{\mu\nu}$

The gravitational action is invariant under general coordinate transformations: gauge invariance



No background to set a scale: background field method

$$g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}$$

The spectrum of the Laplacian of the background metric defines a scale



The gravitational action is invariant under general coordinate transformations: gauge invariance



Introduction of a gauge fixing term:

Faddeev-Popov procedure

$$Z = \int \mathscr{D}h_{\mu\nu} \mathscr{D}\bar{C}_{\alpha} \mathscr{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C]}$$

 $Z = \int \mathscr{D}h_{\mu\nu} \mathscr{D}\bar{C}_{\alpha} \mathscr{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C]}$ Faddeev-Popov ghosts

$$Z = \int \mathscr{D}h_{\mu\nu} \mathscr{D}\bar{C}_{\alpha} \mathscr{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C]}$$

Faddeev-Popov ghosts

$$S_{\rm gf}[\bar{g};h] = \frac{1}{2\alpha} \int d^d x \sqrt{\bar{g}} \, \bar{g}^{\mu\nu} F_{\mu}[\bar{g};h] F_{\nu}[\bar{g};h]$$
$$F_{\mu}[\bar{g};h] = \bar{\nabla}_{\nu} h^{\nu}{}_{\mu} - \frac{1+\beta}{d} \bar{\nabla}_{\mu} h$$
$$S_{\rm gh}[\bar{g};h,\bar{C},C] = \int d^d x \sqrt{\bar{g}} \, \bar{C}_{\alpha} \, \mathscr{M}^{\alpha}{}_{\beta}[\bar{g};h] C^{\beta}$$

$$Z = \int \mathscr{D}h_{\mu\nu} \mathscr{D}\bar{C}_{\alpha} \mathscr{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C]}$$

Faddeev-Popov ghosts

$$S_{\rm gf}[\bar{g};h] = \frac{1}{2\alpha} \int d^d x \sqrt{\bar{g}} \, \bar{g}^{\mu\nu} F_{\mu}[\bar{g};h] F_{\nu}[\bar{g};h]$$
$$F_{\mu}[\bar{g};h] = \bar{\nabla}_{\nu} h^{\nu}{}_{\mu} - \frac{1+\beta}{d} \, \bar{\nabla}_{\mu} h$$
$$S_{\rm gh}[\bar{g};h,\bar{C},C] = \int d^d x \sqrt{\bar{g}} \, \bar{C}_{\alpha} \, \mathscr{M}^{\alpha}{}_{\beta}[\bar{g};h] C^{\beta}$$

Gauge-fixing term breaks split symmetry!

$$Z = \int \mathscr{D}h_{\mu\nu} \mathscr{D}\bar{C}_{\alpha} \mathscr{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C]}$$

Faddeev-Popov ghosts

$$S_{\rm gf}[\bar{g};h] = \frac{1}{2\alpha} \int d^d x \sqrt{\bar{g}} \, \bar{g}^{\mu\nu} F_{\mu}[\bar{g};h] F_{\nu}[\bar{g};h]$$
$$F_{\mu}[\bar{g};h] = \bar{\nabla}_{\nu} h^{\nu}{}_{\mu} - \frac{1+\beta}{d} \bar{\nabla}_{\mu} h$$
$$S_{\rm gh}[\bar{g};h,\bar{C},C] = \int d^d x \sqrt{\bar{g}} \, \bar{C}_{\alpha} \, \mathcal{M}^{\alpha}{}_{\beta}[\bar{g};h] C^{\beta}$$

Gauge-fixing term breaks split symmetry!

Harmless breaking

$$Z = \int \mathscr{D}h_{\mu\nu} \mathscr{D}\bar{C}_{\alpha} \mathscr{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C]}$$

Faddeev-Popov ghosts

$$S_{\rm gf}[\bar{g};h] = \frac{1}{2\alpha} \int d^d x \sqrt{\bar{g}} \, \bar{g}^{\mu\nu} F_{\mu}[\bar{g};h] F_{\nu}[\bar{g};h]$$
$$F_{\mu}[\bar{g};h] = \bar{\nabla}_{\nu} h^{\nu}{}_{\mu} - \frac{1+\beta}{d} \bar{\nabla}_{\mu} h$$
$$S_{\rm gh}[\bar{g};h,\bar{C},C] = \int d^d x \sqrt{\bar{g}} \, \bar{C}_{\alpha} \, \mathscr{M}^{\alpha}{}_{\beta}[\bar{g};h] C^{\beta}$$

Gauge-fixing term breaks split symmetry!

Harmless breaking

Introducing regulators

 $\Delta S_k[\bar{\Phi};\Phi] = \Delta S_k^h[\bar{g};h] + \Delta S_k^{\bar{C}C}[\bar{g};\bar{C},C]$

$$Z = \int \mathscr{D}h_{\mu\nu} \mathscr{D}\bar{C}_{\alpha} \mathscr{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C]}$$

Faddeev-Popov ghosts

$$S_{\rm gf}[\bar{g};h] = \frac{1}{2\alpha} \int d^d x \sqrt{\bar{g}} \, \bar{g}^{\mu\nu} F_{\mu}[\bar{g};h] F_{\nu}[\bar{g};h]$$
$$F_{\mu}[\bar{g};h] = \bar{\nabla}_{\nu} h^{\nu}{}_{\mu} - \frac{1+\beta}{d} \bar{\nabla}_{\mu} h$$
$$S_{\rm gh}[\bar{g};h,\bar{C},C] = \int d^d x \sqrt{\bar{g}} \, \bar{C}_{\alpha} \, \mathscr{M}^{\alpha}{}_{\beta}[\bar{g};h] C^{\beta}$$

Gauge-fixing term breaks split symmetry!

Harmless breaking

Introducing regulators

$$\Delta S_k[\bar{\Phi};\Phi] = \Delta S_k^h[\bar{g};h] + \Delta S_k^{\bar{C}C}[\bar{g};\bar{C},C]$$

$$\Delta S_k^h[\bar{g};h] = \frac{1}{2} \int d^d x \sqrt{\bar{g}} h_{\mu\nu} R_k^{\mu\nu,\alpha\beta}(-\bar{\nabla}^2) h_{\alpha\beta}$$
$$\Delta S_k^h[\bar{g};h] = \int d^d x \sqrt{\bar{g}} \bar{C}_\alpha R_{k,\beta}^\alpha(-\bar{\nabla}^2) C^\beta$$

$$Z = \int \mathscr{D}h_{\mu\nu} \mathscr{D}\bar{C}_{\alpha} \mathscr{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C]}$$

Faddeev-Popov ghosts

$$S_{\rm gf}[\bar{g};h] = \frac{1}{2\alpha} \int d^d x \sqrt{\bar{g}} \, \bar{g}^{\mu\nu} F_{\mu}[\bar{g};h] F_{\nu}[\bar{g};h]$$
$$F_{\mu}[\bar{g};h] = \bar{\nabla}_{\nu} h^{\nu}{}_{\mu} - \frac{1+\beta}{d} \bar{\nabla}_{\mu} h$$
$$S_{\rm gh}[\bar{g};h,\bar{C},C] = \int d^d x \sqrt{\bar{g}} \, \bar{C}_{\alpha} \, \mathscr{M}^{\alpha}{}_{\beta}[\bar{g};h] C^{\beta}$$

Gauge-fixing term breaks split symmetry!

Harmless breaking

Introducing regulators

$$\Delta S_k[\bar{\Phi};\Phi] = \Delta S_k^h[\bar{g};h] + \Delta S_k^{\bar{C}C}[\bar{g};\bar{C},C]$$

Regulators break split symmetry and (quantum) gauge invariance!

$$\Delta S_k^h[\bar{g};h] = \frac{1}{2} \int d^d x \sqrt{\bar{g}} h_{\mu\nu} R_k^{\mu\nu,\alpha\beta}(-\bar{\nabla}^2) h_{\alpha\beta}$$
$$\Delta S_k^h[\bar{g};h] = \int d^d x \sqrt{\bar{g}} \bar{C}_{\alpha} R_{k,\beta}^{\alpha}(-\bar{\nabla}^2) C^{\beta}$$

$$Z_{k}[\mathcal{J}] = \int \mathcal{D}h_{\mu\nu} \mathcal{D}\bar{C}_{\alpha} \mathcal{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C] - \Delta S_{k}[\bar{\Phi};\Phi] + \int d^{d}x \sqrt{\bar{g}} \mathcal{J} \cdot \Phi} \equiv e^{W_{k}[\mathcal{J}]}$$

$$Z_{k}[\mathcal{J}] = \int \mathscr{D}h_{\mu\nu} \mathscr{D}\bar{C}_{\alpha} \mathscr{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C] - \Delta S_{k}[\bar{\Phi};\Phi] + \int d^{d}x \sqrt{\bar{g}} \mathscr{J} \cdot \Phi} \equiv e^{W_{k}[\mathscr{J}]}$$

$$\Gamma_k = \Gamma_k[\bar{\Phi}; \Phi]$$

$$Z_{k}[\mathcal{J}] = \int \mathscr{D}h_{\mu\nu} \mathscr{D}\bar{C}_{\alpha} \mathscr{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C] - \Delta S_{k}[\bar{\Phi};\Phi] + \int d^{d}x \sqrt{\bar{g}} \mathscr{J} \cdot \Phi} \equiv e^{W_{k}[\mathcal{J}]}$$

$$\Gamma_k = \Gamma_k[\bar{\Phi};\Phi]$$

$$\partial_t \Gamma_k[\bar{\Phi}, \Phi] = \frac{1}{2} \operatorname{STr}\left[\left(\Gamma_k^{(0,2)}[\bar{\Phi}, \Phi] + \mathbb{R}_k \right)^{-1} \partial_t \mathbb{R}_k \right]$$

$$Z_{k}[\mathcal{J}] = \int \mathcal{D}h_{\mu\nu} \mathcal{D}\bar{C}_{\alpha} \mathcal{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C] - \Delta S_{k}[\bar{\Phi};\Phi] + \int d^{d}x \sqrt{\bar{g}} \mathcal{J} \cdot \Phi} \equiv e^{W_{k}[\mathcal{J}]}$$

$$\Gamma_k = \Gamma_k[\bar{\Phi}; \Phi]$$

$$\partial_t \Gamma_k[\bar{\Phi}, \Phi] = \frac{1}{2} \operatorname{STr} \left[\left(\Gamma_k^{(0,2)}[\bar{\Phi}, \Phi] + \mathbb{R}_k \right)^{-1} \partial_t \mathbb{R}_k \right]$$

The effective average action is a functional of two fields;

Integrating the flow and taking *k*=0 leads to an effective average action that depends on two fields, but background independence is guaranteed by BRST symmetry;

$$Z_{k}[\mathcal{J}] = \int \mathscr{D}h_{\mu\nu} \mathscr{D}\bar{C}_{\alpha} \mathscr{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C] - \Delta S_{k}[\bar{\Phi};\Phi] + \int d^{d}x \sqrt{\bar{g}} \mathscr{J} \cdot \Phi} \equiv e^{W_{k}[\mathscr{J}]}$$

$$\Gamma_k = \Gamma_k[\bar{\Phi}; \Phi]$$

$$\partial_t \Gamma_k[\bar{\Phi}, \Phi] = \frac{1}{2} \operatorname{STr} \left[\left(\Gamma_k^{(0,2)}[\bar{\Phi}, \Phi] + \mathbb{R}_k \right)^{-1} \partial_t \mathbb{R}_k \right]$$

The effective average action is a functional of two fields;

Integrating the flow and taking *k*=0 leads to an effective average action that depends on two fields, but background independence is guaranteed by BRST symmetry;


$$Z_{k}[\mathcal{J}] = \int \mathcal{D}h_{\mu\nu} \mathcal{D}\bar{C}_{\alpha} \mathcal{D}C^{\beta} e^{-S[\bar{g}+h] - S_{gf}[\bar{g};h] - S_{gh}[\bar{g};h,\bar{C},C] - \Delta S_{k}[\bar{\Phi};\Phi] + \int d^{d}x \sqrt{\bar{g}}\mathcal{J} \cdot \Phi} \equiv e^{W_{k}[\mathcal{J}]}$$

In complete analogy: construction of effective average action

$$\Gamma_k = \Gamma_k[\bar{\Phi}; \Phi]$$

$$\partial_t \Gamma_k[\bar{\Phi}, \Phi] = \frac{1}{2} \operatorname{STr} \left[\left(\Gamma_k^{(0,2)}[\bar{\Phi}, \Phi] + \mathbb{R}_k \right)^{-1} \partial_t \mathbb{R}_k \right]$$

The effective average action is a functional of two fields;

Integrating the flow and taking *k*=0 leads to an effective average action that depends on two fields, but background independence is guaranteed by BRST symmetry;



very little is known about the Gribov problem in quantum gravity and how it can affect background independence



Our starting point was a path integral over Riemannian metrics

Our starting point was a path integral over Riemannian metrics

However...

Our starting point was a path integral over Riemannian metrics

However...

 $h_{\mu
u}$ can fluctuate widely

Our starting point was a path integral over Riemannian metrics

However...

 $h_{\mu
u}\,$ can fluctuate widely

metric can be degenerate



metric can change signature

Our starting point was a path integral over Riemannian metrics



Such a linear split of the metric might introduce many spurious configurations in the nonperturbative realm!

Our starting point was a path integral over Riemannian metrics



Such a linear split of the metric might introduce many spurious configurations in the nonperturbative realm!

Alternative:

avoid the previous problems + cover the space of Riemannian metrics

Our starting point was a path integral over Riemannian metrics



Such a linear split of the metric might introduce many spurious configurations in the nonperturbative realm!

Alternative:

$$g_{\mu\nu} = \bar{g}_{\mu\alpha} \left(\mathrm{e}^{\bar{g}^{-1}h} \right)^{\alpha}_{\nu}$$

avoid the previous problems + cover the space of Riemannian metrics

In the path integral, should we adopt different variables that lead to the same field equations in the case of GR?

Our starting point was a path integral over Riemannian metrics



Such a linear split of the metric might introduce many spurious configurations in the nonperturbative realm!

Alternative:

 $g_{\mu\nu} =$

$$\bar{g}_{\mu\alpha}\left(\mathrm{e}^{\bar{g}^{-1}h}\right)^{\alpha}_{\nu}$$

avoid the previous problems + cover the space of Riemannian metrics

In the path integral, should we adopt different variables that lead to the same field equations in the case of GR?

Palatini
$$(g_{\mu\nu}, \Gamma^{\alpha}_{\beta\sigma})$$
 or pure e^{a}_{μ} or $(e^{a}_{\mu}, \omega^{bc}_{\nu})$

Our starting point was a path integral over Riemannian metrics



Such a linear split of the metric might introduce many spurious configurations in the nonperturbative realm!

Alternative:

avoid the previous problems + cover the space of Riemannian metrics

In the path integral, should we adopt different variables that lead to the same field equations in the case of GR?

Palatini
$$(g_{\mu\nu}, \Gamma^{\alpha}_{\beta\sigma})$$
 or pure e^{a}_{μ} or $(e^{a}_{\mu}, \omega^{bc}_{\nu})$

No a priori reason to choose one formulation instead of the other

$$h_{\mu\nu} = h_{\mu\nu}^{\rm TT} + \bar{\nabla}_{\mu}\xi_{\nu} + \bar{\nabla}_{\nu}\xi_{\mu} + \bar{\nabla}_{\mu}\bar{\nabla}_{\nu}\sigma - \frac{1}{d}\bar{g}_{\mu\nu}\bar{\nabla}^{2}\sigma + \frac{1}{d}\bar{g}_{\mu\nu}h$$

$$h_{\mu\nu} = h_{\mu\nu}^{\rm TT} + \bar{\nabla}_{\mu}\xi_{\nu} + \bar{\nabla}_{\nu}\xi_{\mu} + \bar{\nabla}_{\mu}\bar{\nabla}_{\nu}\sigma - \frac{1}{d}\bar{g}_{\mu\nu}\bar{\nabla}^{2}\sigma + \frac{1}{d}\bar{g}_{\mu\nu}h$$

In the cosmological term:
$$\frac{2}{16\pi G_k} \int d^d x \sqrt{g} \Lambda_k$$

$$h_{\mu\nu} = h_{\mu\nu}^{\rm TT} + \bar{\nabla}_{\mu}\xi_{\nu} + \bar{\nabla}_{\nu}\xi_{\mu} + \bar{\nabla}_{\mu}\bar{\nabla}_{\nu}\sigma - \frac{1}{d}\bar{g}_{\mu\nu}\bar{\nabla}^{2}\sigma + \frac{1}{d}\bar{g}_{\mu\nu}h$$

In the cosmological term:
$$\frac{2}{16\pi G_k} \int d^d x \sqrt{g} \Lambda_k$$

Making use of the exponential parametrization:

$$\frac{2}{16\pi G_k} \int \mathrm{d}^d x \sqrt{\bar{g}} \,\mathrm{e}^{\frac{1}{2}\mathrm{Tr}\,h} \Lambda_k$$

$$h_{\mu\nu} = h_{\mu\nu}^{\rm TT} + \bar{\nabla}_{\mu}\xi_{\nu} + \bar{\nabla}_{\nu}\xi_{\mu} + \bar{\nabla}_{\mu}\bar{\nabla}_{\nu}\sigma - \frac{1}{d}\bar{g}_{\mu\nu}\bar{\nabla}^{2}\sigma + \frac{1}{d}\bar{g}_{\mu\nu}h$$

In the cosmological term: $\frac{2}{16\pi G_k}\int d^dx \sqrt{g} \Lambda_k$

Making use of the exponential parametrization:

$$\frac{2}{16\pi G_k} \int \mathrm{d}^d x \sqrt{\bar{g}} \,\mathrm{e}^{\frac{1}{2}\mathrm{Tr}\,h} \Lambda_k$$

Just the trace of h couples to the cosmological constant

$$h_{\mu\nu} = h_{\mu\nu}^{\rm TT} + \bar{\nabla}_{\mu}\xi_{\nu} + \bar{\nabla}_{\nu}\xi_{\mu} + \bar{\nabla}_{\mu}\bar{\nabla}_{\nu}\sigma - \frac{1}{d}\bar{g}_{\mu\nu}\bar{\nabla}^{2}\sigma + \frac{1}{d}\bar{g}_{\mu\nu}h$$

In the cosmological term: $\frac{2}{16\pi G_k} \int d^d x \sqrt{g} \Lambda_k$

Making use of the exponential parametrization:

$$\frac{2}{16\pi G_k} \int \mathrm{d}^d x \sqrt{\bar{g}} \,\mathrm{e}^{\frac{1}{2}\mathrm{Tr}\,h} \Lambda_k$$

Just the trace of *h* couples to the cosmological constant

Can we eliminate the trace of *h* by a suitable choice of gauge? YES!

$$h_{\mu\nu} = h_{\mu\nu}^{\rm TT} + \bar{\nabla}_{\mu}\xi_{\nu} + \bar{\nabla}_{\nu}\xi_{\mu} + \bar{\nabla}_{\mu}\bar{\nabla}_{\nu}\sigma - \frac{1}{d}\bar{g}_{\mu\nu}\bar{\nabla}^{2}\sigma + \frac{1}{d}\bar{g}_{\mu\nu}h$$

In the cosmological term: $\frac{2}{16\pi G_k} \int d^d x \sqrt{g} \Lambda_k$

Making use of the exponential parametrization:

$$\frac{2}{16\pi G_k} \int \mathrm{d}^d x \sqrt{\bar{g}} \mathrm{e}^{\frac{1}{2}\mathrm{Tr}\,h} \Lambda_k$$

Just the trace of h couples to the cosmological constant

Can we eliminate the trace of *h* by a suitable choice of gauge? YES!

Recall

$$F_{\mu}[\bar{g};h] = \bar{\nabla}_{\nu}h^{\nu}{}_{\mu} - \frac{1+eta}{d}\bar{\nabla}_{\mu}h \qquad \text{w/} \qquad \beta \to -\infty \qquad \text{freezes the fluctuations of trace of }h$$

$$h_{\mu\nu} = h_{\mu\nu}^{\rm TT} + \bar{\nabla}_{\mu}\xi_{\nu} + \bar{\nabla}_{\nu}\xi_{\mu} + \bar{\nabla}_{\mu}\bar{\nabla}_{\nu}\sigma - \frac{1}{d}\bar{g}_{\mu\nu}\bar{\nabla}^{2}\sigma + \frac{1}{d}\bar{g}_{\mu\nu}h$$

In the cosmological term: $\frac{2}{16\pi G_k}\int d^dx \sqrt{g} \Lambda_k$

Making use of the exponential parametrization:

$$\frac{2}{16\pi G_k} \int \mathrm{d}^d x \sqrt{\bar{g}} \,\mathrm{e}^{\frac{1}{2}\mathrm{Tr}\,h} \Lambda_k$$

Just the trace of h couples to the cosmological constant

Can we eliminate the trace of *h* by a suitable choice of gauge? YES!

Recall

$$F_{\mu}[\bar{g};h] = \bar{\nabla}_{\nu} h^{\nu}{}_{\mu} - \frac{1+\beta}{d} \bar{\nabla}_{\mu} h \qquad \mathbf{w} / \qquad \beta \to -\infty$$

$$\frac{2}{16\pi G_k} \int \mathrm{d}^d x \sqrt{\bar{g}} \,\mathrm{e}^{\frac{1}{2}\mathrm{Tr}\,h} \Lambda_k \to \frac{2}{16\pi G_k} \int \mathrm{d}^d x \sqrt{\bar{g}} \,\Lambda_k$$

freezes the fluctuations of trace of h

field independent: does not contribute to the flow

From the previous discussion, we see that the cosmological constant drops out Beta functions of other couplings will not depend on the cosmological constant The cosmological constant still runs but even its flow does not depend on the cc itself

From the previous discussion, we see that the cosmological constant drops out Beta functions of other couplings will not depend on the cosmological constant The cosmological constant still runs but even its flow does not depend on the cc itself

Moreover

 $\det(g_{\mu\nu}) = \det(\bar{g}_{\mu\alpha}) e^{\frac{1}{2} \operatorname{Tr} h} \quad \text{with } \operatorname{Tr} h = 0 \quad \det(g_{\mu\nu}) = \det(\bar{g}_{\mu\alpha})$

From the previous discussion, we see that the cosmological constant drops out Beta functions of other couplings will not depend on the cosmological constant The cosmological constant still runs but even its flow does not depend on the cc itself

Moreover

$$\det(g_{\mu\nu}) = \det(\bar{g}_{\mu\alpha}) e^{\frac{1}{2} \operatorname{Tr} h} \quad \text{with } \operatorname{Tr} h = 0 \quad \det(g_{\mu\nu}) = \det(\bar{g}_{\mu\alpha})$$

The determinant of the metric (to be path-integrated) is fixed and equals the determinant of the background metric

From the previous discussion, we see that the cosmological constant drops out Beta functions of other couplings will not depend on the cosmological constant The cosmological constant still runs but even its flow does not depend on the cc itself

Moreover

$$\det(g_{\mu\nu}) = \det(\bar{g}_{\mu\alpha}) e^{\frac{1}{2} \operatorname{Tr} h} \quad \text{with } \operatorname{Tr} h = 0 \quad \det(g_{\mu\nu}) = \det(\bar{g}_{\mu\alpha})$$



From the previous discussion, we see that the cosmological constant drops out Beta functions of other couplings will not depend on the cosmological constant The cosmological constant still runs but even its flow does not depend on the cc itself

Moreover

$$\det(g_{\mu\nu}) = \det(\bar{g}_{\mu\alpha}) e^{\frac{1}{2}\operatorname{Tr} h} \quad \text{with } \operatorname{Tr} h = 0 \quad \det(g_{\mu\nu}) = \det(\bar{g}_{\mu\alpha}) e^{\frac{1}{2}\operatorname{Tr} h}$$



The existence of the UV-fixed point imposes severe constraints on the RG-flow

Quantum Scale Symmetry

The existence of the UV-fixed point imposes severe constraints on the RG-flow



Finitely many free parameters

The existence of the UV-fixed point imposes severe constraints on the RG-flow



The existence of the UV-fixed point imposes severe constraints on the RG-flow



The existence of the UV-fixed point imposes severe constraints on the RG-flow



This is a hint for a fixed point that is not deeply non-perturbative

The existence of the UV-fixed point imposes severe constraints on the RG-flow



This is a hint for a fixed point that is not deeply non-perturbative

$$G_k = k^{2-d}g_k$$
 $\beta = (d-2)g_k + F_k(g)$ $d = 4$ $\beta = 2g_k + F_k(g)$

The existence of the UV-fixed point imposes severe constraints on the RG-flow



This is a hint for a fixed point that is not deeply non-perturbative

$$G_k = k^{2-d}g_k$$
 $\beta = (d-2)g_k + F_k(g)$ $d = 4$ $\beta = 2g_k + F_k(g)$

Perturbative calculations: asymptotic freedom

$$d = 2$$
 $\beta = F_k(g)$

The existence of the UV-fixed point imposes severe constraints on the RG-flow



Weinberg proposes the asymptotic safety scenario in a special edition in celebration to Einstein's centenary

Weinberg proposes the asymptotic safety scenario in a special edition in celebration to Einstein's centenary

No technique to test such a conjecture in four dimensions

Weinberg proposes the asymptotic safety scenario in a special edition in celebration to Einstein's centenary

No technique to test such a conjecture in four dimensions

M. Reuter seminal work on FRG & QG

1996-1998 arXiv:hep-th/9605030 Weinberg proposes the asymptotic safety scenario in a special edition in celebration to Einstein's centenary

No technique to test such a conjecture in four dimensions



1996-1998 arXiv:hep-th/9605030


Weinberg proposes the asymptotic safety scenario in a special edition in celebration to Einstein's centenary

No technique to test such a conjecture in four dimensions



In order to extract the effects of quantum-gravity fluctuations, one has to solve the "fully quantum equations of motion"

In order to extract the effects of quantum-gravity fluctuations, one has to solve the "fully quantum equations of motion"

Interplay between Quantum gravity and Particle Physics

Cosmology and black holes

In order to extract the effects of quantum-gravity fluctuations, one has to solve the "fully quantum equations of motion"

Interplay between Quantum gravity and Particle Physics

Cosmology and black holes



Quite active research topic in the field

In order to extract the effects of quantum-gravity fluctuations, one has to solve the "fully quantum equations of motion"



In order to extract the effects of quantum-gravity fluctuations, one has to solve the "fully quantum equations of motion"



It seems that we cannot exclude the possibility of quantum gravity to be described by an asymptotically safe standard QFT

It seems that we cannot exclude the possibility of quantum gravity to be described by an asymptotically safe standard QFT

Non-perturbative tools are mandatory in this case - Big Challenge!

It seems that we cannot exclude the possibility of quantum gravity to be described by an asymptotically safe standard QFT

Non-perturbative tools are mandatory in this case - Big Challenge!

Besides technicalities, there are important conceptual open problems:

It seems that we cannot exclude the possibility of quantum gravity to be described by an asymptotically safe standard QFT

Non-perturbative tools are mandatory in this case - Big Challenge!

Besides technicalities, there are important conceptual open problems:

How to transport everything that we have learnt so far to the Lorentzian setting?

It seems that we cannot exclude the possibility of quantum gravity to be described by an asymptotically safe standard QFT

Non-perturbative tools are mandatory in this case - Big Challenge!

Besides technicalities, there are important conceptual open problems:

How to transport everything that we have learnt so far to the Lorentzian setting?

Is the theory unitary?

It seems that we cannot exclude the possibility of quantum gravity to be described by an asymptotically safe standard QFT

Non-perturbative tools are mandatory in this case - Big Challenge!

Besides technicalities, there are important conceptual open problems:

How to transport everything that we have learnt so far to the Lorentzian setting?

Is the theory unitary?

Do we have a complete RG-trajectory that emanates from the UV to our IR?

It seems that we cannot exclude the possibility of quantum gravity to be described by an asymptotically safe standard QFT

Non-perturbative tools are mandatory in this case - Big Challenge!

Besides technicalities, there are important conceptual open problems:

How to transport everything that we have learnt so far to the Lorentzian setting?

Is the theory unitary?

Do we have a complete RG-trajectory that emanates from the UV to our IR?

How to connect the results obtained with the FRG and other non-perturbative schemes?

Stay tuned or join us! :)